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Jonathan Ziveyi 1, Craig Blackburn 2, Michael Sherris 3

1 School of Risk and Actuarial Studies, University of New South Wales, Sydney, Australia, 2052

2 School of Risk and Actuarial Studies, CEPAR, University of New South Wales, Sydney, Australia, 2052

3 School of Risk and Actuarial Studies. Michael Sherris is a Chief Investigator at the ARC Centre of Excellence in Population Ageing Research, University of New South Wales, Sydney, Australia, 2052

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Pricing European Options on Deferred Insurance Contracts

Jonathan Ziveyi^a, Craig Blackburn^b, Michael Sherris^c

^aSchool of Risk and Actuarial Studies, University of New South Wales, Sydney, Australia, 2052 ^bSchool of Risk and Actuarial Studies, CEPAR, University of New South Wales, Sydney, Australia, 2052 ^cSchool of Risk and Actuarial Studies, CEPAR, University of New South Wales, Sydney, Australia, 2052

Abstract

This paper considers the pricing of European call options written on pure endowment and deferred life annuity contracts, also known as guaranteed annuity options. These contracts provide a guarantee value at maturity of the option. The contract valuation is dependent on stochastic interest rate and mortality processes. We assume single-factor stochastic squareroot processes for both interest rate and mortality intensity, with mortality being a timeinhomogeneous process. We then derive the pricing partial differential equation (PDE) and the corresponding transition density PDE for options written on deferred contracts. The general solution of the pricing PDE is derived as a function of the transition density function. We solve the transition density PDE by first transforming it to a system of characteristic PDEs using Laplace transform techniques and then applying the method of characteristics. Once an explicit expression for the density function is found, we then use sparse grid quadrature techniques to generate European call option prices on deferred insurance products. This approach can easily be generalised to other contracts which are driven by similar stochastic processes presented in this paper. We test the sensitivity of the option prices by varying independent parameters in our model. As option maturity increases, the corresponding option prices significantly increase. The effect of miss-pricing the guaranteed annuity value is analysed, as is the benefit of replacing the whole-life annuity with a term annuity to remove volatility of the old age population.

JEL Classification C63, G22, G13, G12

Keywords: Mortality risk, Deferred insurance products, European options, Laplace Transforms

Email addresses: j.ziveyiQunsw.edu.au (Jonathan Ziveyi), c.blackburnQunsw.edu.au (Craig Blackburn), m.sherrisQunsw.edu.au (Michael Sherris)

1. Introduction

Insurance companies and annuity providers are increasingly exposed to the risk of ever improving mortality trends across all ages, with a greater portion of survivors living beyond 100 years (Carriere [9], Currie et al. [12] and CMI(2005) among others). Such mortality improvements, coupled with the unavailability of suitable hedging instruments pose significant challenges to annuity providers as this results in longer periods of annuity payments than initial forecasts. At present mortality risk is non-tradable (Blake et al. [7]) and there is no market to hedge these risks other than reinsurance. Blake and Burrows [6] highlight that annuity providers have been trying to hedge mortality risk using costly means such as the construction of hedged portfolios of long-term bonds (with no mortality risk). There is need for an active longevity market where mortality risk can easily be priced, traded and hedged. A number of securities that can make up this market have already been proposed in the literature and these include longevity bonds, mortality derivatives, securitised products among others (Bauer [2] and Blake et al. [7]).

Blake et al. [7], Bauer [1], and Bauer [2] demonstrate that if mortality risk can be traded through securities such as longevity bonds and swaps then the techniques developed in financial markets can be adapted and implemented for mortality risk. A number of papers have also developed models for pricing guaranteed annuity options. Milevsky and Promislow [21] develop algorithms for valuing mortality contingent claims by taking the underlying securities as defaultable coupon paying bonds with the time of death as a stopping time. Boyle and Hardy [8] use the numeraire approach to value options written on guaranteed annuities. They detail the challenges experienced in the UK where long-dated and low guaranteed rates were provided relative to the high prevailing interest rates in the 1970's. Interest rates fell significantly in the 1990's leading to sharp increases in the value of guaranteed contracts and this had a significant impact on annuity providers' profitability.

A significant number of empirical studies have been presented showing that mortality trends are generally improving and the future development of mortality rates are considered stochastic. Amongst proposed stochastic models include Milevsky and Promislow [21], Dahl et al. [13], Biffis [4]. A number of stochastic mortality models have been strongly inspired by interest rate term-structure modelling literature (Cox et al. [11], Dahl et al. [13], Litterman and Scheinkman [19]) as well as stochastic volatility models such as that proposed in Heston [17].

The main aim of this paper is to devise a novel approach for the pricing and hedging of deferred mortality contingent claims with special emphasis on pure endowment options and deferred immediate annuity options. We start off by devising techniques for pricing deferred insurance contracts, which are the underlying assets. Analytical solutions can be derived for pure endowment contracts using the forward measure approach, however, this is not possible for deferred immediate annuities where analytical approximation techniques have mostly been used as in Singleton and Umantsev [23] when valuing options on coupon paying bonds. Having devised models for the underlying securities, we then outline techniques for

valuing European style options written on these contracts.

In developing our framework, we assume that the interest rate dynamics is driven by a singlefactor stochastic square-root process while the time-inhomogeneous mortality dynamics is a one-factor version of the model proposed in Biffis [5]. The long-term mean reversion level of the mortality process is a time-varying function following the Weibull mortality law. This works as a reference mortality level for each age in the cohort. The model definition guarantees positive mortality rates, although mortality intensity can fall below our reference rate, careful selection of the parameters limits this occurring.

We use hedging arguments and Ito's Lemma to derive partial differential equation (PDE) for options written on the deferred insurance contracts. We also present the backward Kolmogorov PDE satisfied by the two stochastic processes under consideration. Most of the techniques we develop in this paper are drawn from in Chiarella and Ziveyi [10] where the dynamics of the underlying security evolve under the influence of stochastic volatility. We present the general solution of the pricing PDE by using Duhamel's principle. This solution is a function of the joint probability function which is the solution of the Kolmogorov PDE. We solve the transition density PDE with the aide of Laplace transform techniques thereby obtaining an explicit expression for the joint transition density function. Using the explicit density function, we then use sparse grid quadrature methods to price options on deferred insurance contracts.

The remainder of this paper is organised as follows. Section 2 presents the modelling framework for the interest rate and mortality rate processes. We then provide the option pricing framework in Section 3. It is in this section where we derive the option pricing PDE and the corresponding backward Kolmogorov PDE for the density function. With the general solution of the pricing PDE presented, we then outline a step-by-step approach for solving the transition density PDE using Laplace transform techniques. The explicit expressions for the deferred pure endowment and deferred annuity contracts together with their corresponding option prices are presented in Section 4. All numerical results are presented in Section 5. Section 6 concludes the paper. Where appropriate, the long derivations and proofs are included to the appendices.

2. Modelling Framework

The intensity based modelling of credit risky securities has a number of parallels with mortality modelling (Lando [18] and Biffis [5]). We are interested in the first stopping time, τ , of the intensity process $\mu(t; x)$, for a person aged x at time zero. Starting with a filtered probability space $(\Omega, \mathcal{F}, \mathbb{F}, \mathbb{P})$, where \mathbb{P} is the real-world probability measure. The information at time-t is given by $\mathbb{F} = \mathbb{G} \vee \mathbb{H}$. The sub-filtration \mathbb{G} contains all financial and actuarial information except the actual time of death. There are two \mathbb{G} -adapted short-rate processes, r(t) and $\mu(t; x)$; representing the instantaneous interest and mortality processes respectively. The sub-filtration \mathbb{H} is the σ -algebra with death information. Let $N(t) := \mathbb{1}_{\tau \leq t}$ be an indicator function, if the compensator $A(t) = \int_0^t \mu(s; x) ds$ is a predictable process of N(t) then dM(t) = dN(t) - dA(t) is a \mathbb{P} -martingale, where $dA(t) = \mu(t; x) dt$. There also exists another measure where dM(t) is a \mathbb{Q} -martingale, under which the compensator becomes $dA(t) = \mu^Q(t; x) dt$, with $\mu^Q(t; x) = (1 + \phi(t))\mu(t; x)$ and $\phi(t) \geq -1$. The function $\phi(t)$ may represent unsystematic risk of the insurance contract, we assume $\phi(t) = 0$ in this paper.

Proposition 2.1. In the absence of arbitrage opportunities there exists an equivalent martingale measure \mathbb{Q} where $C(t,T,r,\mu;x)$ is the t-value of an option contract with a pay-off function, $P(T,r,\mu;x)$, at time-T. The payment of $P(T,r,\mu;x)$, which is a \mathbb{G} -adapted process, is conditional on survival to the start of the period T, otherwise the value is zero. The time t-value of the option can be represented as

$$C(t,T,r,\mu;x) = \mathbb{E}^{Q} \left[e^{-\int_{t}^{T} r(u)du} P(T,r,\mu;x) \mathbb{1}_{\tau > T} |\mathcal{F}_{t} \right]$$
$$= \mathbb{1}_{\tau > t} \mathbb{E}^{Q} \left[e^{-\int_{t}^{T} [r(u)+\mu(u;x)]du} P(T,r,\mu;x) |\mathcal{G}_{t} \right]$$
(2.1)

Proof: The law of iterated expectations can be used to show this, see Bielecki and Rutkowski [3] or Biffis [5] for detail. \Box

In our framework, time-T is always the option maturity age. If $P(T, r, \mu; x) \equiv 1$ then our contract resembles a credit risky zero coupon bond. In actuarial terms, this is a pure endowment contract written at time-t that receives 1 at time T if the holder is still alive. If $P(T, r, \mu; x)$ is the value of a stream of payments starting at T, conditional on survival, then we are pricing a deferred immediate annuity. The contract value, $P(T, r, \mu; x)$, can also take the form of an option pay-off. In this scenario the strike price, K, represents a guaranteed value at time-T on an endowment or annuity contract.

One approach to solving equation (2.1) is to use a forward measure approach. If we use $P(T, r, \mu; x)$ as numeraire, we can rewrite (2.1) as

$$C(t, T, r, \mu; x) = \mathbb{1}_{\tau > t} P(t, r, \mu; x) \mathbb{E}^{Q^{P}} \left[P(T, r, \mu; x) | \mathcal{F}_{t} \right]$$
(2.2)

where \mathbb{E}^{Q^P} is our new forward probability measure. When $P(T, r, \mu; x)$ is the value of a general payment stream, closed form solutions to this problem do not exist. One solution is to use Monte-Carlo simulations or numerical approximation methods. These approximation methods are derived by Singleton and Umantsev [23] for coupon bearing bond options in a general affine framework, while Schrager [22] proposes a numerical approximation method for pricing guaranteed annuity options in a Gaussian affine framework. For a guaranteed annuity option, this requires using the deferred annuity as a numeraire as outlined in Boyle and Hardy [8].

In our approach we derive a closed-form joint density function for interest and mortality rates. This method uses Laplace transform techniques allowing us to directly solve equation (2.1) under the risk neutral measure. With a closed-form density function we use numerical integration to price various types of insurance contracts. We use sparse grids quadrature techniques to evaluate double integral expressions.

2.1. Interest Rate Model

In our filtered probability space $(\Omega, \mathcal{F}, \mathbb{F}, \mathbb{P})$, defined above, the G-adapted short-rate process r, represents the instantaneous interest rate on risk-free securities. There exists an equivalent martingale measure \mathbb{Q} where the arbitrage-free price of a risk-free zero coupon bond is given by

$$B(t,T) = \mathbb{E}^{Q} \left[e^{-\int_{t}^{T} r(u)du} |\mathcal{G}_{t} \right], \qquad (2.3)$$

where T is the maturity of the bond and r(u) is the short-rate process. We model the interest rate as a single factor affine process (Duffie and Kan [16] and Dai and Singleton [15]). The short-rate is modelled as a Cox-Ingersoll-Ross (CIR) process with the risk-neutral dynamics defined as

$$dr(t) = \kappa_r(\theta_r - r(t))dt + \sigma_r\sqrt{r(t)}dW^Q, \qquad (2.4)$$

where κ_r , is the speed of mean reversion of r with θ_r being the corresponding long-run mean. The volatility of the process is denoted by σ_r . For the process (2.4) to be guaranteed positive, Cox et al. [11] show the parameters must satisfy the following condition $\frac{2\kappa_r}{\theta_r} > \sigma_r^2$. The explicit solution of equation (2.4) can be represented as

$$B(t,T) = e^{\alpha_r(t,T) - \beta_r(t,T)r(t)}$$

$$(2.5)$$

where α_r and β_r are expressions of the form

$$\begin{split} \alpha_r(t,T) = & \frac{2\kappa_r \theta_r}{\sigma_r^2} \log\left[\frac{2\gamma_r e^{\frac{(\kappa_r + \gamma_r)(T-t)}{2}}}{(\gamma_r + \kappa_r)(e^{\gamma_r(T-t)} - 1) + 2\gamma_r}\right]\\ \beta_r(t,T) = & \frac{2(e^{\gamma_r(T-t)} - 1)}{(\gamma_r + \kappa_r)(e^{\gamma_r(T-t)} - 1) + 2\gamma_r}\\ \gamma_r = & \sqrt{\kappa_r^2 + 2\sigma_r^2} \end{split}$$

2.2. Mortality Model

We model mortality as an affine process. The choice of this process is problematic, since Luciano and Vigna [20] show that a time-homogeneous model like equation (2.4) does not represent mortality rates very well. Another approach is to use time-inhomogeneous models as proposed in Biffis [5] and Dahl and Moller [14]. In this paper we use a 1-factor version of the Biffis [5] 2-factor square-root diffusion model. The mortality intensity process is modelled

$$d\mu(t;x) = \kappa_{\mu}(m(t) - \mu(t;x))dt + \sigma_{\mu}\sqrt{\mu(t;x)}dW(t)$$
(2.6)

where

 $\sigma_{\mu} = \Sigma_{\mu} \sqrt{m(t)}$

By using similar arguments as in Cox et al. [11], the mortality process is positive definite if parameters are chosen such that $2\kappa_{\mu} > \Sigma_{\mu}^2$. Biffis [5] chooses m(t) to be a deterministic function given by

$$m(x+t) = \frac{c}{\theta^{c}}(x+t)^{c-1}$$
(2.7)

which is the Weibull mortality law; we adopt this functional form as it fits well to observed mortality trends.

Using the same filtered probability space $(\Omega, \mathcal{F}, \mathbb{F}, \mathbb{P})$, the G-adapted process μ , represents the instantaneous mortality rate. There exists an equivalent martingale measure \mathbb{Q} such that the survival probability can be represented as

$$S(t,T;x) = \mathbb{E}^{Q} \left[e^{-\int_{t}^{T} \mu(u)du} |\mathcal{G}_{t} \right]$$
(2.8)

The general solution of equation (2.8) can be shown to be

$$S(t,T;x) = e^{\alpha_{\mu}(t,T;x) - \beta_{\mu}(t,T;x)\mu(x,t)}$$
(2.9)

where $\alpha_{\mu}(t,T;x)$ and $\beta_{\mu}(t,T;x)$ are the solutions to the following ordinary differential equations

$$\frac{\partial}{\partial t}\beta_{\mu}(t,T;x) = 1 - \kappa_{\mu}(t;x)\beta_{\mu}(t,T;x) - \frac{1}{2}\left(\Sigma_{\mu}(x,t)\right)^{2}\left(\beta_{\mu}(t,T;x)\right)^{2}m(t),$$
(2.10)

$$\frac{\partial}{\partial t}\alpha_{\mu}(t,T;x) = -\kappa_{\mu}(t;x)m(t)(x,t)\beta_{\mu}(t,T;x), \qquad (2.11)$$

with $\beta_{\mu}(T,T;x) = 0$ and $\alpha_{\mu}(T,T;x) = 0$, see Duffie and Kan [16]. A closed form solution of $\beta_{\mu}(t,T;x)$ and $\alpha_{\mu}(t,T;x)$ does not exist for this model, but can be solved via numerical techniques such as the Runge-Kutta methods.

3. The Option Pricing Framework

By using hedging arguments and Ito's Lemma the time-t value of an option, $C(t, T, r, \mu; x)$, written on an insurance contract with value $P(T, r, \mu; x)$ at time-T, is the solution to the partial differential equation (PDE)

$$\frac{\partial C}{\partial t}(t,T,r,\mu;x) = \mathcal{L}C(t,T,r,\mu;s) - r_x C(t,T,r,\mu;s), \qquad (3.1)$$

where

$$\mathcal{L} = \kappa_r (\theta_r - r) \frac{\partial}{\partial r} + \kappa_\mu (m(t) - \mu) \frac{\partial}{\partial \mu} + \frac{1}{2} \sigma_r^2 r \frac{\partial^2}{\partial r^2} + \frac{1}{2} \sigma_\mu^2 \mu \frac{\partial^2}{\partial \mu^2}, \qquad (3.2)$$

is the Dynkin operator and $0 < r, \mu < \infty$, $r_x = r + \mu$ and t < T. Detailed discussion on deriving PDEs like (3.1) can be found in Chiarella and Ziveyi [10]. Equation (3.1) is solved subject to a boundary condition; which is the payoff of the option at the time it matures. Also associated with the system of SDEs (2.4) and (2.6) is the transition density function, which we denote here as $G(\psi, r, \mu; r_T, \mu_T, x)$ with $\psi = T - t$ being the time-to-maturity. Chiarella and Ziveyi (2011) show that SDEs like (2.4) and (2.6) satisfy the associated Kolmogorov backward PDE

$$\frac{\partial G}{\partial \psi}(\psi, r, \mu; r_0, \mu_0, x) = \mathcal{L}G(\psi, r, \mu; r_0, \mu_0, x).$$
(3.3)

Equation (3.3) is solved subject to the boundary condition

$$G(T, r, \mu; r_T, \mu_T, x) = \delta(r - r_T)\delta(\mu - \mu_T)$$
(3.4)

where $\delta(\cdot)$ is the Dirac Delta function, r_T is the instantaneous short-rate while μ_T is the instantaneous mortality rate.

3.1. General Solution of the Option Pricing Problem

By using the law of iterated expectation as detailed in Biffis [5] and the independence assumption between r and μ , equation (2.1) can be re-expressed as

$$C(t,T,r,\mu;x) = \mathbb{1}_{\tau > t} \mathbb{E}^{Q} \left[e^{-\int_{t}^{T} [r(u) + \mu(u;x)] du} \mathbb{E}^{Q} \left[P(T,r,\mu;x) | \mathcal{G}_{T} \right] | \mathcal{G}_{t} \right].$$
(3.5)

The problem then becomes that of finding the contract value, $P(T, r, \mu; x)$, at time-T discounted to time-t; we present the general solution in the proposition below.

Proposition 3.1. By making use of Duhamel's principle, the general solution of equation (3.5) can be represented as

$$C(t,T,r,\mu;x) = 1_{\tau > t} \mathbb{E}^{Q} \left[e^{-\int_{t}^{T} [r(u) + \mu(u;x)] du} |\mathcal{G}_{t} \right] \int_{0}^{\infty} \int_{0}^{\infty} P(T,r,\mu;x) G(T,r,\mu;w_{1},w_{2}) dw_{1} dw_{2}$$

$$= 1_{\tau > t} e^{\alpha_{r}(t,T) - \beta_{r}(t,T)r(t)} \times e^{\alpha_{\mu}(t,T;x) - \beta_{\mu}(t,T;x)\mu(x,t)} \times \int_{0}^{\infty} \int_{0}^{\infty} P(T,r,\mu;x) G(T,r,\mu;w_{1},w_{2}) dw_{1} dw_{2}, \qquad (3.6)$$

where $G(T, r, \mu; w_1, w_2)$ is the solution of the transition density PDE (3.3).

Proof: A detailed proof of Duhamel's principle is presented in Chiarella and Ziveyi [10]. \Box

Now that we have managed to present the general solution of our pricing function, we only need to solve (3.3) for the density function which is the only unknown term in equation (3.6). We accomplish this in the next section.

3.2. Applying the Laplace Transform

In solving (3.3) we make use of Laplace transform techniques to transform the PDE to a corresponding system of characteristic PDEs; we present this in the next proposition.

Proposition 3.2. Laplace transform of equation (3.3) can be represented as

$$\frac{\partial \tilde{G}}{\partial T} + \left\{ \frac{1}{2} \sigma_r^2 s_r^2 - \kappa_r s_r \right\} \frac{\partial \tilde{G}}{\partial s_r} + \left\{ \frac{1}{2} \sigma_\mu^2 s_\mu^2 - \kappa_\mu s_\mu \right\} \frac{\partial \tilde{G}}{\partial s_\mu} \\
= \left[(\kappa_r \theta_r - \sigma_r^2) s_r + (\kappa_\mu m(t) - \sigma_\mu^2) s_\mu + \kappa_r + \kappa_\mu \right] \tilde{G} + f_1(\tau, s_\mu) + f_2(\tau, s_r)$$
(3.7)

where

$$f_1(\tau, s_\mu) = \left(\frac{1}{2}\sigma_r^2 - \kappa_r \theta_r\right) \tilde{G}(\tau, 0, s_\mu)$$
(3.8)

$$f_2(\tau, s_r) = \left(\frac{1}{2}\sigma_{\mu}^2 - \kappa_{\mu}m(t)\right)\tilde{G}(\tau, s_r, 0).$$
(3.9)

Equation (3.7) is to be solved subject to the initial condition

$$\tilde{G}(0, r, \mu; r_0, \mu_0, x) = e^{-s_r r_0 - s_\mu \mu_0}$$
(3.10)

Proof: Refer to Appendix 1.

Proposition 3.3. The solution to equation (3.3) can be represented as

$$\tilde{G}(T, s_r, s_\mu; r_0, \mu_0, x) = \left(\frac{2\kappa_r}{\sigma_r^2 s_r (e^{\kappa_r T} - 1) + 2\kappa_r}\right)^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \times \left(\frac{2\kappa_\mu}{\sigma_\mu^2 s_\mu (e^{\kappa_\mu T} - 1) + 2\kappa_\mu}\right)^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_r^2}} \\
exp\left[\{\kappa_r + \kappa_\mu\}T\right] \times exp\left\{-\frac{2s_r \kappa_r e^{\kappa_r T}}{\sigma_r^2 s(e^{\kappa_r T} - 1) + 2\kappa_r}r_0 - \frac{2s_\mu \kappa_\mu e^{\kappa_\mu T}}{\sigma_\mu^2 s_\mu (e^{\kappa_\mu T} - 1) + 2\kappa_\mu}\mu_0\right\} \\
\times \left[\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1; \frac{4\kappa_r^2 e^{\kappa_r T}}{(e^{\kappa_r T} - 1)(\sigma_r^2 s(e^{\kappa_r T} - 1) + 2\kappa_r)}r_0\right) \\
+ \Gamma\left(\frac{2\Phi_\mu}{\sigma_\mu^2} - 1; \frac{4\kappa_\mu^2 e^{\kappa_\mu T}}{(e^{\kappa_\mu T} - 1)(\sigma_\mu^2 s_\mu (e^{\kappa_\mu T} - 1) + 2\kappa_\mu)}\mu_0\right) - 1\right]$$
(3.11)

where

$$\Phi_r = \kappa_r \theta_r$$
$$\Phi_\mu = \kappa_\mu m(t)$$

Proof: Refer to Appendix 2.

Proposition 3.4. The inverse Laplace transform of equation (3.11) can be represented as

$$G(T, r, \mu; r_0, \mu_0, x) = exp \left\{ -\frac{2\kappa_r}{\sigma_r^2 (e^{\kappa_r T} - 1)} (r_0 e^{\kappa_r T} + r) - \frac{2\kappa_\mu}{\sigma_\mu^2 (e^{\kappa_\mu T} - 1)} (\mu_0 e^{\kappa_\mu T} + \mu) \right\} \\ \times \left(\frac{r_0 e^{\kappa_r T}}{r} \right)^{\frac{\Phi_r}{\sigma_r^2} - \frac{1}{2}} \times \left(\frac{\mu_0 e^{\kappa_\mu T}}{\mu} \right)^{\frac{\Phi_\mu}{\sigma_\mu^2} - \frac{1}{2}} I_{\frac{2\Phi_r}{\sigma_r^2} - 1} \left(\frac{4\kappa_r}{\sigma_r^2 (e^{\kappa_r T} - 1)} (r \times r_0 e^{\kappa_r T})^{\frac{1}{2}} \right) \\ \times I_{\frac{2\Phi_\mu}{\sigma_\mu^2} - 1} \left(\frac{4\kappa_\mu}{\sigma_\mu^2 (e^{\kappa_\mu T} - 1)} (\mu \times \mu_0 e^{\kappa_\mu T})^{\frac{1}{2}} \right) \times \frac{2\kappa_r e^{\kappa_r T}}{\sigma_r^2 (e^{\kappa_r T} - 1)} \frac{2\kappa_\mu e^{\kappa_\mu T}}{\sigma_\mu^2 (e^{\kappa_\mu T} - 1)}$$
(3.12)

Proof: Refer to Appendix 3.

4. Deferred Insurance Contracts

When pricing deferred insurance contracts, we first define general functions to denote such contracts. The value of a pure endowment contract at option maturity is given by

$$F_E(T, r, \mu; x) = \mathbb{E}^Q \left[e^{-\int_T^{T_m} [r(u) + \mu(u)] du} \Big| \mathcal{G}_T \right], \qquad (4.1)$$

where T_m is the time when the pure endowment contract matures and T is the option maturity date. Similarly the payoff for a deferred immediate annuity, when the annuity starts at the option maturity date, T, can be represented as

$$F_A(T, r, \mu; x) = \sum_{i=T}^{\omega} \mathbb{E}^Q \left[e^{-\int_T^i [r(u) + \mu(u)] du} \Big| \mathcal{G}_T \right].$$

$$(4.2)$$

More complicated structures can arise when an option on a deferred immediate annuity expires before the annuity start time such that

$$F_A(T,r,\mu;x) = \mathbb{E}^Q \left[e^{-\int_T^{T+h} [r(u)+\mu(u)] du} \left[\sum_{i=T+h}^{\omega} \mathbb{E}^Q \left[e^{-\int_{T+h}^i [r(u)+\mu(u)] du} \Big| \mathcal{G}_{T+h} \right] \right] \Big| \mathcal{G}_T \right], \quad (4.3)$$

where the filtration, \mathcal{G}_T , defines the option maturity date, ω is the maximum age in the cohort for a whole life immediate annuity and h is the time between option maturity and annuity start age. Equations like (4.3) naturally lead to the pricing of American style options on life

insurance contracts; we have reserved this for future research. In this paper, we are going to focus on functional forms in equations (4.1) and (4.2).

Given our affine definition of the short-rate interest rate and mortality processes we have an explicit solution for the expectation in equation (4.1) given by

$$F_E(T, r, \mu; \varphi_1, \varphi_2, x) = e^{\alpha_r (T_m - T) - \beta_r (T_m - T)\varphi_1} \times e^{\alpha_\mu (T_m - T; x) - \beta_\mu (T_m - T; x)\varphi_2}, \qquad (4.4)$$

and for equation (4.2) we have

$$F_A(T, r, \mu; \varphi_1, \varphi_2, x) = \sum_{i=T}^{\omega} \left[e^{\alpha_r (i-T) - \beta_r (i-T)\varphi_1} \times e^{\alpha_\mu (i-T; x) - \beta_\mu (i-T; x)\varphi_2} \right].$$
 (4.5)

For a pure endowment contract, when $T = T_m$ the option and the contract matures at the same time such that

$$F_E(T, r, \mu; \varphi_1, \varphi_2, x) = 1,$$
(4.6)

and by definition of a probability density function we obtain

$$\int_0^\infty \int_0^\infty G(T, r, \mu; \varphi_1, \varphi_2, x) d\varphi_1 d\varphi_2 = 1, \qquad (4.7)$$

implying that the price of a pure endowment option is simply

$$C(t, T, r, \mu; x) = 1_{\tau > t} e^{\alpha_r (T-t) - \beta_r (T-t)r(t)} e^{\alpha_\mu (T-t;x) - \beta_\mu (T-t;x)\mu(t;x)},$$
(4.8)

which is known at time-t.

4.1. Options on Deferred Insurance Contracts

We take the perspective of the insurer and focus on European call options written on insurance contracts defined in equations (4.1) and (4.2). An option on a pure endowment contract has a pay-off at option maturity which we represent here as

$$P(T, r, \mu; x) = \max\left[0, \mathbb{E}^{Q}\left[e^{-\int_{T}^{T_{m}} [r(u)+\mu(u)]du} \middle| \mathcal{G}_{T}\right] - K\right],$$

$$(4.9)$$

where T_m is the time when the pure endowment contract matures, T is the option maturity date and K is the guaranteed value. Similarly the pay-off of a deferred immediate annuity, when the annuity starts at the same time as the option expires, time-T, can be represented as

$$P(T, r, \mu; x) = \max\left[0, \sum_{i=T}^{\omega} \mathbb{E}^{Q}\left[e^{-\int_{T}^{i} [r(u)+\mu(u)]du} \middle| \mathcal{G}_{T}\right] - K\right],$$
(4.10)

where K is the guaranteed annuity value at time-T.

By substituting equations (4.9) and (4.10) into equation (3.6) we obtain the option price of a pure endowment and deferred immediate annuity, respectively. Such option prices can be represented as

$$C(t, T, r, \mu; x) = 1_{\tau > t} e^{\alpha_r (T-t) - \beta_r (T-t)r(t)} e^{\alpha_\mu (T-t;x) - \beta_\mu (T-t;x)\mu(t;x)} \\ \times \int_0^\infty \int_0^\infty \max \left[0, e^{\alpha_r (T_m - T) - \beta_r (T_m - T)\varphi_1} \times e^{\alpha_\mu (T_m - T;x) - \beta_\mu (T-T;x)\varphi_2} - K \right] G(T, r, \mu; \varphi_1, \varphi_2, x) d\varphi_1 d\varphi_2 .$$

$$(4.11)$$

and

$$C(t,T,r,\mu;x) = 1_{\tau > t} e^{\alpha_r (T-t) - \beta_r (T-t)r(t)} e^{\alpha_\mu (T-t;x) - \beta_\mu (T-t;x)\mu(t;x)} \\ \times \int_0^\infty \int_0^\infty \max \left[0, \sum_{i=T}^\omega e^{\alpha_r (i-T) - \beta_r (i-T)\varphi_1} \times e^{\alpha_\mu (i-T;x) - \beta_\mu (i-T;x)\varphi_2} - K \right] G(T,r,\mu;\varphi_1,\varphi_2,x) d\varphi_1 d\varphi_2$$
(4.12)

5. Numerical Results

5.1. Model Parameters

Figure 5.1 shows the joint density function for varying values of κ_{μ} and Σ_{μ} for a cohort aged 65. The interest rate parameters are fixed, typical parameter values are given in 1. We note different shapes of the joint density function when we vary the the speed of reversion parameter, κ_{μ} and volatility Σ_{μ} in the mortality process, equation 2.6. The effect of increasing κ_{μ} increases the peak mortality density, hence a lower survival probability. Increasing the volatility, Σ_{μ} , significantly increases the mortality intensity dispersion. Practically, the integral limits to infinity are not required in equations (4.11) and (4.12). We can see that the interest rate and mortality intensity processes decay to zero, and we will use these observed limits when performing numerical integration.

Figure 5.2 shows the mortality intensity probability density function at various ages for a volatility fixed at $\Sigma_{\mu} = 0.15$; in this figure we have integrated the interest rate process. The figure is truncated at a mortality intensity of 0.12 to observe detail at the younger ages, all the intensity distributions decay to zero and integrate to one to indicate a proper probability distribution. As the cohort's age increases, i.e. time to maturity, the peak of the mortality intensity is increasing. The dispersion of mortality intensity is increasing with age, giving a high level of uncertainty of mortality intensity as our cohort ages.

We price guaranteed pure endowment and deferred annuity options in the framework proposed derived in section 3. Table 1 shows the parameters of the Weibull function used in this paper as a base level of our cohort at age 50, equation (2.7). These are the initial values used



Figure 5.1: Probability Density Function - Age = 65



Figure 5.2: Mortality Intensity for Different Ages

by Biffis [5]. Table 1 also contains the parameter values of the interest rate and mortality stochastic processes derived in sections 2.1 and 2.2 respectively. Unless otherwise specified these values are used in the analysis that follows.

κ_r	$ heta_r$	σ_r	Σ_{μ}	c	θ	r_0	μ_0
0.05	0.05	0.01	0.15	10.841	86.165	0.03	m(0)

Table 1: Stochastic process and Weibull parameters

In this analysis we model a cohort aged 50 at time zero and pricing options maturing at time T. The time-to-maturity is defined as $\psi = T - t$.

To derive the survival probabilities at time zero we solve the system of ordinary differential equations given in equation (2.10). The shape of the survival curves can be controlled by changing the κ_{μ} and Σ_{μ} parameters of the mortality stochastic process. Figures 5.3a and 5.3c show the survival curves at time zero for various values of κ_{μ} and Σ_{μ} respectively. Lower values of κ_{μ} , corresponding to a slower speed of mean-reversion, produces survival curves



Figure 5.3: Survival Probability and Force of Mortality

with higher survival probabilities for a fixed Σ_{μ} . Similarly, we fix κ_{μ} and observe the survival curves as Σ_{μ} varies. Higher values of Σ_{μ} increase the survival probability in the older ages.

The force of mortality for our cohort aged 50 at time zero is defined as

$$\mu(t,T;x) = -\frac{d}{dT} \log \left[\mathbb{E}^Q \left[e^{-\int_t^T \mu(u) du} \Big| \mathcal{G}_t \right] \right]$$
$$= -\frac{d}{dT} \left[\alpha_\mu(T-t;x) - \beta_\mu(T-t;x)\mu(t;x) \right].$$
(5.1)

The force of mortality for varying levels of κ_{μ} and Σ_{μ} is shown in Figures 5.3b and 5.3d respectively. The scale of each figure changes for ages below and above age 90. In Figure 5.3b changing values of κ_{μ} produces a large variety of curves below age 90; for a fixed Σ_{μ} , above age 90 the force of mortality for each κ_{μ} does not diverge significantly. The opposite can be seen in Figure 5.3d, for a fixed κ_{μ} , varying Σ_{μ} produces a divergence in the force of mortality over age 90. By varying these two parameters we have a flexible mortality model that can cover a large variety of survival curves.



Figure 5.4: Probability Density Function

We also wish to observe the probability density function of the mortality process at some time-T in the future. This can be done by integrating the joint density function, equation (3.12), with respect to r_0 . Figure 5.4 shows the mortality probability density function with T = 65 for varying values of κ_{μ} and Σ_{μ} , interest rate process parameters are given in Table 1. The peak in the density function in Figure 5.4a corresponds to the lowest value of κ_{μ} . The peak density decreases and mortality intensity dispersion increases as κ_{μ} increases. Figure 5.4b shows the mortality probability density function as Σ_{μ} varies. The peak in the density function is greatest at the lowest volatility level. Increasing Σ_{μ} decreases (and then increases) density level while the mortality intensity dispersion increases.

5.2. Option Pricing

Using the parameters from Table 1 we can analyse the effect of a changing κ_{μ} on pure endowment and deferred annuity contract valuations and option prices. For the pure endowment contract, a payment of one is made to the contract holder at the maturity of contract if they are alive. Similarly for the deferred annuity contract holder, a periodic payment of one is made to the contract holder while they are still alive, the annuity is a whole-life immediate type. The guaranteed value in the options, K, will typically be given as the model market value contract valuation at time-T. All the options we are pricing are European call options based on a cohort aged 50 at time zero.

5.2.1. Pure Endowment Pricing

Figure 5.5 shows results for a pure endowment with an option maturing at T = 65 and the contract maturity varying from age 65 to 100. When the option maturity and contract maturity are both 65, the contract valuation at time-T is 1 and the option price is zero; similar to payoff of an at-the-money option. As the maturity of the pure endowment contract increases to 100 the effect of varying the κ_{μ} parameter can be seen in Figure 5.5a. The contracts with lower κ_{μ} 's have higher values at time-T. For the option price, the guaranteed value is the model market value of the pure endowment contract valued at time-T. The pure endowment contract value at time-t is given in Figure 5.5b.

Figure 5.5c shows the option prices for a the pure endowment with guarantee value given in Figure 5.5a. Even though the contract values are always decreasing with age, the option prices are increasing until age 75 or 80. Then the prices of the options decrease with age, this is due to the shape of the survival curves in Figure 5.3a. The guaranteed values are dependent on the model market value of the pure endowment contract at time T. In reality, an insurer may prefer to analyse option prices for a fixed guaranteed value and varying assumptions in κ_{μ} , and we perform this miss-pricing analysis below.

The option prices relative to the face-value of the pure endowment valuation at time-t are shown in Figure 5.5d. These percentages are increasing in age but flattens out after 90. This shows that the longer the contract maturity date relative to the option maturity date, the higher the cost of the hedge. The option price as a percentage of the contract face value is relatively insensitive to changes in κ_{μ} in this example.



Figure 5.5: Pure Endowment Options

Unfortunately, the guaranteed values at time-T are dependent on our choice of κ_{μ} . Since

$\begin{array}{c} \kappa_{\mu} \\ \text{Age} \end{array}$	0.015	0.025	0.035	0.045	0.055	0.065	\bar{K}
70	0.8022	0.7986	0.7955	0.7928	0.7905	0.7885	0.7947
75	0.6321	0.6231	0.6153	0.6086	0.6029	0.5979	0.6133
80	0.4830	0.4650	0.4501	0.4376	0.4271	0.4181	0.4468
85	0.3487	0.3193	0.2963	0.2781	0.2634	0.2516	0.2929
90	0.2287	0.1899	0.1627	0.1431	0.1285	0.1175	0.1617
95	0.1301	0.0912	0.0681	0.0536	0.0440	0.0374	0.0707
100	0.0622	0.0334	0.0201	0.0132	0.0093	0.0070	0.0242

Table 2: Guaranteed Value for different κ_{μ}

κ_{μ} Age	0.015	0.025	0.035	0.045	0.055	0.065	$\begin{array}{c} \kappa_{\mu} \\ \text{Age} \end{array}$	0.015	0.025	0.035	0.045	0.055	0.065
70	0.0062	0.0043	0.0030	0.0020	0.0014	0.0009	70	1.10	0.77	0.53	0.37	0.25	0.17
75	0.0139	0.0083	0.0045	0.0022	0.0009	0.0004	75	3.19	1.92	1.04	0.50	0.21	0.08
80	0.0258	0.0135	0.0050	0.0011	0.0001	0.0000	80	8.11	4.25	1.59	0.34	0.03	0.00
85	0.0397	0.0188	0.0043	0.0001	0.0000	0.0000	85	19.06	9.04	2.06	0.05	0.00	0.00
90	0.0476	0.0200	0.0021	0.0000	0.0000	0.0000	90	41.39	17.42	1.85	0.00	0.00	0.00
95	0.0422	0.0145	0.0001	0.0000	0.0000	0.0000	95	84.00	28.98	0.29	0.00	0.00	0.00
100	0.0270	0.0065	0.0000	0.0000	0.0000	0.0000	100	156.86	38.15	0.00	0.00	0.00	0.00
		(a) C	ption Pric	e	(b) Percentage								

 Table 3: Pure Endowment Options

we cannot know the correct κ_{μ} , we analyse miss-pricing of the options by fixing guaranteed values. This time we choose a guarantee value that is the mean of the contract values at time-T, denoted by \bar{K} , see Table 2. As the contract maturity increases, there are larger changes in model market value of the contracts. Table 3 shows the option prices and the percentages of face-values for guarantee values of \bar{K} given in Table 2. Under these assumptions the option prices and percentages of face-value are highly dependent on κ_{μ} . This shows that in reality offering guaranteed pure endowments contracts that mature over age 80 may be extremely expensive.

5.2.2. Deferred Immediate Annuity Pricing

The pricing of deferred immediate annuity options is similar to the previous sections. From equation (4.2) we can see that the annuity payment is the sum of pure endowment contracts as shown in equation (4.1). In this section the annuity contract is a whole life deferred immediate annuity, with the annuity contract deferred to time-T.

Figure 5.6 shows our analysis for a changing option maturity date, T, varying from 55 to 100. Figure 5.6a shows the guarantee value or model market value of the deferred immediate annuity at time-T. As expected, the lower values of κ_{μ} have higher guarantee values. The contract values at time-t for different annuity start year, time-T, are shown in Figure 5.6b; with the corresponding option prices in Figure 5.6c. We can see a relative low percentage of option price to contract face-value when the deferment period is short, see Figure 5.6d.

This shows a relative increase in the cost of options with deferment period, there is a quite large difference in the percentage for varying κ_{μ} at the longer deferment periods. Unlike pure endowment option prices, the relative prices of the options on deferred immediate annuities are sensitive to the value of κ_{μ} .

κ_{μ} Age	0.015	0.025	0.035	0.045	0.055	0.065	\bar{K}
55	20.9465	20.1327	19.6371	19.2976	19.0496	18.8612	19.6541
60	18.3791	17.5232	16.9917	16.6207	16.3445	16.1307	16.9983
65	16.5308	15.5718	14.9750	14.5567	14.2447	14.0029	14.9803
70	14.9051	13.8099	13.1300	12.6581	12.3098	12.0425	13.1426
75	13.2778	12.0286	11.2739	10.7643	10.3973	10.1214	11.3106
80	11.5325	10.1404	9.3374	8.8174	8.4571	8.1954	9.4134
85	9.6840	8.2205	7.4280	6.9410	6.6177	6.3906	7.5470
90	7.8290	6.3911	5.6667	5.2460	4.9790	4.7983	5.8183
95	6.1091	4.7963	4.1838	3.8483	3.6450	3.5126	4.3492
100	4.6438	3.5298	3.0478	2.7977	2.6523	2.5606	3.2053

Table 4: Guaranteed Value for different κ_{μ}

Similar to the previous section, we also wish to test the sensitivity of the option prices for fixed guarantee values, rather than the model market values. The last column of Table 4 gives the average 'fair-values' of the annuity contracts at time-T. Table 5 shows the option prices and percentage of option price to contract value. Since these contract valuations depend on the remaining life of the contract holder, we see higher sensitivity to the guarantee value than pure endowment contracts.

κ_{μ} Age	0.015	0.025	0.035	0.045	0.055	0.065	κ_{μ} Age	0.015	0.025	0.035	0.045	0.055	0.065
55	1.1978	0.4545	0.1068	0.0139	0.0012	0.0001	55	6.58	2.50	0.59	0.08	0.01	0.00
60	1.1419	0.4420	0.0954	0.0081	0.0003	0.0000	60	8.14	3.15	0.68	0.06	0.00	0.00
65	1.0991	0.4236	0.0782	0.0029	0.0000	0.0000	65	10.32	3.99	0.74	0.03	0.00	0.00
70	1.0466	0.4052	0.0853	0.0035	0.0000	0.0000	70	13.43	5.25	1.11	0.05	0.00	0.00
75	0.9425	0.3650	0.0978	0.0113	0.0002	0.0000	75	17.50	6.96	1.91	0.22	0.00	0.00
80	0.7848	0.3041	0.1000	0.0236	0.0031	0.0001	80	22.97	9.46	3.26	0.80	0.11	0.01
85	0.5696	0.2112	0.0764	0.0249	0.0068	0.0014	85	29.58	12.45	4.97	1.75	0.51	0.11
90	0.3455	0.1151	0.0418	0.0156	0.0058	0.0020	90	37.23	15.83	6.93	2.98	1.22	0.46
95	0.1680	0.0460	0.0153	0.0057	0.0023	0.0010	95	45.02	19.05	8.80	4.24	2.08	1.01
100	0.0643	0.0130	0.0036	0.0012	0.0005	0.0002	100	51.34	21.34	10.13	5.21	2.81	1.56
		(a) (Option Prio	ce		(b) Percentage							

Table 5: Deferred Whole-Life Annuity

5.3. Volatilities of Interest and Mortality Processes

We have two stochastic processes affecting the price of our deferred insurance contracts. Interest rate processes are generally known as highly volatile in the short term, with the volatility of prices decreasing with time as the process reverts to its long term mean value. The effect of changing interest rate volatility is presented in Figure 5.7a. We can see the



Figure 5.6: Deferred Immediate Annuity Options

difference in relative price in the contracts with the shortest deferment period. This difference is due to interest rate volatility converging around the age of 75 for our selected parameters.



(a) Annuity - Changing Interest Rate Volatility (b) Annuity - Changing Mortality Volatility

Figure 5.7: Volatility

The situation is reversed for mortality volatility. The different levels of volatility significantly change the mortality rates at the older ages, see Figure 5.3d. We present the effect of varying volatility on relative option prices in Figure 5.7b. These effects are small at the younger ages, but as the deferment age increases the relative price of the options starts to diverge. Prolonged deferment periods may not be practical, a realistic deferment period for a person aged 50 may be between 15 and 30 years. At a 15 year deferment the effect of volatility is small on the relative option price. The difference at age 80 is between 1 and 4 percent.

5.4. Term Annuities

In this analysis we compare the differences in the option price for a deferred whole life annuity and a deferred term annuity contract. The deferment period for both contracts is 15 years; corresponding to a person currently aged 50 with a guaranteed annuity value at aged 65. The guarantee value is the model market value annuity price at age 65. The term annuity contracts matures at age 80 and removes the risk associated with the old age uncertainty in the cohort.

Table 6 shows the guarantee value of each contract with varying values of Σ_{μ} and κ_{μ} . As expected, deferred whole life annuity contracts are more expensive compared to the corresponding deferred term annuity contracts. These findings also hold for the corresponding option prices which are higher for whole life contracts as showing in Table 7. The percentage of option price to contract value is shown in Table 8. We can see the percentage for wholelife annuities show fairly large changes in κ_{μ} for a given Σ_{μ} . Whereas for the term annuity options, the percentage is relatively insensitive to κ_{μ} for a given Σ_{μ} .

κ_{μ} Σ_{μ}	0.015	0.025	0.035	0.045	0.055	0.065	κ_{μ} Σ_{μ}	0.015	0.025	0.035	0.045	0.055	0.065
0.03	11.1076	11.0146	10.9335	10.8601	10.7896	10.7562	0.03	16.1282	15.3085	14.7803	14.3988	14.1015	13.9045
0.06	11.1066	11.0098	10.9278	10.8576	10.7975	10.7454	0.06	16.1811	15.3370	14.7984	14.4158	14.1285	13.9046
0.09	11.1071	11.0104	10.9284	10.8583	10.7981	10.7461	0.09	16.2698	15.3954	14.8418	14.4504	14.1569	13.9286
0.12	11.1077	11.0112	10.9293	10.8593	10.7991	10.7470	0.12	16.3875	15.4747	14.9013	14.4978	14.1961	13.9618
0.15	11.1111	11.0120	10.9300	10.8600	10.7997	10.7476	0.15	16.5308	15.5718	14.9750	14.5567	14.2447	14.0029
	(a) l	Deferred Ter	m Annuity	Contract Pr	rice		(b) Def	erred Whole	e-Life Annui	ty Contract	Price		

 Table 6: Contract Prices

κ_{μ} Σ_{μ}	0.015	0.025	0.035	0.045	0.055	0.065	κ_{μ} Σ_{μ}	0.015	0.025	0.035	0.045	0.055	0.065
0.03	0.0314	0.0296	0.0287	0.0296	0.0330	0.0260	0.03	0.0692	0.0606	0.0559	0.0548	0.0575	0.0470
0.06	0.0333	0.0330	0.0325	0.0320	0.0315	0.0312	0.06	0.0742	0.0668	0.0622	0.0588	0.0562	0.0545
0.09	0.0355	0.0352	0.0349	0.0342	0.0337	0.0332	0.09	0.0803	0.0723	0.0672	0.0631	0.0601	0.0579
0.12	0.0381	0.0381	0.0378	0.0370	0.0364	0.0357	0.12	0.0873	0.0789	0.0733	0.0686	0.0651	0.0623
0.15	0.0400	0.0414	0.0412	0.0404	0.0397	0.0389	0.15	0.0930	0.0861	0.0801	0.0749	0.0710	0.0676
	(a) De	oforrod Tor	m Annuity	Option P	rico			(b) Defer	red Whole	Life Appr	uity Option	Price	

(a) Deferred Term Annuity Option Price

(b) Deferred Whole-Life Annuity Option Price

Table 7: Option Prices

κ_{μ} Σ_{μ}	0.015	0.025	0.035	0.045	0.055	0.065	κ_{μ} Σ_{μ}	0.015	0.025	0.035	0.045	0.055	0.065
0.03	0.3983	0.3796	0.3719	0.3870	0.4352	0.3439	0.03	0.6035	0.5584	0.5356	0.5401	0.5794	0.4816
0.06	0.4221	0.4226	0.4213	0.4177	0.4150	0.4133	0.06	0.6451	0.6152	0.5951	0.5786	0.5657	0.5582
0.09	0.4494	0.4519	0.4516	0.4466	0.4433	0.4400	0.09	0.6944	0.6632	0.6409	0.6199	0.6040	0.5928
0.12	0.4829	0.4884	0.4891	0.4837	0.4789	0.4735	0.12	0.7499	0.7199	0.6960	0.6712	0.6517	0.6358
0.15	0.5064	0.5304	0.5331	0.5282	0.5227	0.5158	0.15	0.7920	0.7804	0.7571	0.7301	0.7083	0.6877
	(a) Defer	red Term	Annuity O	ption Perce	entage			(b) Defe	erred Whol	e-Life Ann	uity Perce	ntage	

Table 8: Percentage (Option Price / Price)

5.5. Variable Strike

In this section we analyse how the option price changes when we vary the strike price. As the strike goes to zero the price of the option converges to the market value of the contract. Figure 5.8 shows the option price as the strike price increases to the market value. We vary the strike for a whole life deferred annuity contract with option maturities at age 65 and 80. Figures 5.8a and 5.8b show large shifts in the option price when κ_{μ} varies. Smaller shifts in the option price occur when Σ_{μ} is varied.



Figure 5.8: Deferred Whole Life Annuity Option Price

6. Conclusion

We have presented a mathematical framework for pricing deferred insurance contracts and options written on such contracts. The approach involves the derivation of the pricing partial differential equation (PDE) for options written on deferred insurance products and the corresponding Kolmogorov PDE for the joint transition density function. The general solution of the pricing PDE has been presented with the aide of Duhamel's principle and this is a function of the joint transition density, where the transition density is a solution of the Kolmogorov PDE for the two stochastic processes under consideration.

We have outlined a systematic approach for solving the transition density PDE with the aide of Laplace transform techniques. Application of Laplace transforms to the PDE transforms it to a corresponding system of characteristic PDEs which can then be solved by the method of characteristics. This yields an explicit closed-form expression for the bivariate transition density function. The closed-form expression allows us to price option contracts on deferred insurance products without the need to perform Monte Carlo simulations or any numerical approximation techniques. Guided by the shape of the joint transition density function, we establish finite upper bounds for the integrals that appear in the pricing expressions. We then use sparse grids quadrature techniques to solve option pricing functions. Numerical analysis has been provided for the contract values, option prices and sensitivity analysis for varying volatility and interest rates. We have analysed the effects of changing the speed of mean-reversion, κ_{μ} , for the mortality process on contracts and option prices where we noted that lower speeds of mean-reversion will always yield higher contract prices and corresponding higher option prices. We have shown that the guarantee value is dependent on κ_{μ} and the problem of miss-pricing contracts can lead to very high option prices relative to contract value.

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Appendix 1. Proof of Proposition 3.2

By using the Laplace transform definition

$$\mathcal{L}[G] = \int_0^\infty \int_0^\infty e^{-s_r r - s_\mu \mu} G(\psi, r, \mu) dr d\mu \equiv \tilde{G}(\psi, s_r, s_\mu), \tag{A1.1}$$

we can derive the following transforms

$$\begin{split} \mathcal{L}\left[\frac{\partial G}{\partial r}\right] &= -\tilde{G}(\psi, 0, s_{\mu}) + s\tilde{G}(\psi, s_{r}, s_{\mu}), \\ \mathcal{L}\left[\frac{\partial G}{\partial \mu}\right] &= -\tilde{G}(\psi, s_{r}, 0) + s_{\mu}\tilde{G}(\psi, s_{r}, s_{mu}), \\ \mathcal{L}\left[r\frac{\partial G}{\partial \mu}\right] &= -\tilde{G}(\psi, s_{r}, s_{mu}) - s_{r}\frac{\partial \tilde{G}}{\partial s_{r}}(\psi, s_{r}, s_{\mu}), \\ \mathcal{L}\left[\mu\frac{\partial G}{\partial \mu}\right] &= -\tilde{G}(\psi, s_{r}, s_{\mu}) - s_{\mu}\frac{\partial \tilde{G}}{\partial s_{\mu}}(\psi, s_{r}, s_{\mu}), \\ \mathcal{L}\left[r\frac{\partial^{2}G}{\partial r^{2}}\right] &= \tilde{G}(\psi, 0, s_{\mu}) - 2s_{r}\tilde{G}(\psi, s_{r}, s_{\mu}) - s_{r}^{2}\frac{\partial \tilde{G}}{\partial s_{r}}(\psi, s_{r}, s_{\mu}), \\ \mathcal{L}\left[\mu\frac{\partial^{2}G}{\partial \mu^{2}}\right] &= \tilde{G}(\psi, s_{r}, 0) - 2s_{\mu}\tilde{G}(\psi, s_{r}, s_{\mu}) - s_{\mu}^{2}\frac{\partial \tilde{G}}{\partial s_{\mu}}(\psi, s_{r}, s_{\mu}). \end{split}$$

Applying these transforms to equation (3.3) we obtain the results in equation (3.7) of Proposition 3.2.

Appendix 2. Proof of Proposition 3.3

Using the method of characteristics to solve (3.7) with initial condition (3.10). Equation (3.7) can be re-expressed in characteristic form as

$$d\psi = \frac{ds_r}{\frac{1}{2}\sigma_r^2 s_r^2 - \kappa_r s_r} = \frac{ds_\mu}{\frac{1}{2}\sigma_\mu^2 s_\mu^2 - \kappa_\mu s_\mu} = d\tilde{G} \Big/ \Big[\left\{ (\kappa_r \theta_r - \sigma_r^2) s + (\kappa_\mu m(t) - \sigma_\mu^2) s_\mu + \kappa_r + \kappa_\mu \right\} \tilde{G} + f_1(\psi, s_\mu) + f_2(\psi, s_r) \Big]$$
(A2.1)

We solve the first characteristic pair of (A2.1) by integrating to obtain

$$\int d\psi = \frac{2}{\sigma_r^2} \int \frac{ds_r}{s\left(s_r - \frac{2\kappa_r}{\sigma_r^2}\right)}$$

This implies that

$$\psi + c_1 = \frac{1}{\kappa_r} \int \left[\frac{1}{s_r - \frac{2\kappa_r}{\sigma_r^2}} - \frac{1}{s_r} \right] ds_r$$

integrating the RHS gives

$$\kappa_r \psi + c_2 = \ln\left[rac{s_r - rac{2\kappa_r}{\sigma_r^2}}{s_r}
ight]$$

which implies

$$e^{\kappa_r\psi+c_2} = \frac{s_r - \frac{2\kappa_r}{\sigma_r^2}}{s_r}$$

hence

$$c_3 = e^{c_2} = \frac{(s_r - \frac{2\kappa_r}{\sigma_r^2})e^{-\kappa_r\psi}}{s_r}$$
$$= \frac{(\sigma_r^2 s - 2\kappa_r)e^{-\kappa_r\psi}}{\sigma_r^2 s_r}$$

This can be rearranged to solve for s_r

$$c_{3}\sigma_{r}^{2}s = (\sigma_{r}^{2}s - 2\kappa_{r})e^{-\kappa_{r}\psi}$$

$$c_{3}\sigma_{r}^{2}s_{r} - \sigma_{r}^{2}s_{r}e^{-\kappa_{r}\psi} = -2\kappa_{r}e^{-\kappa_{r}\psi}$$

$$s_{r} = \frac{-2\kappa_{r}e^{-\kappa_{r}\psi}}{\sigma_{r}^{2}(c_{3} - e^{-\kappa_{r}\psi})}$$
(A2.2)

Similarly the second characteristic pair of equation (A2.1) can be given by

$$d_3 = \frac{(s_\mu - \frac{2\kappa_\mu}{\sigma_\mu^2})e^{-\kappa_\mu\psi}}{s_\mu}$$
$$= \frac{(\sigma_\mu^2 s_\mu - 2\kappa_\mu)e^{-\kappa_\mu\psi}}{\sigma_\mu^2 s_\mu}$$

with

$$s_{\mu} = \frac{-2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}{\sigma_{\mu}^2(d_3 - e^{-\kappa_{\mu}\psi})}$$
(A2.3)

(A2.4)

The last characteristic pair can be represented as

$$\frac{d\tilde{G}}{d\psi} + \left\{ (\kappa_r \theta_r - \sigma_r^2) s + \left(\kappa_\mu m(t) - \sigma_\mu^2 \right) s_\mu + \kappa_r + \kappa_\mu \right\} \tilde{G} = f_1(\psi, s_\mu) + f_2(\psi, s_r) \quad (A2.5)$$

The integrating factor of equation (A2.5) is

$$R(\psi) = \exp\left[\int \left\{ (\sigma_r^2 - \kappa_r \theta_r) s_r + \left(\sigma_\mu^2 - \kappa_\mu m(t)\right) s_\mu - \kappa_r - \kappa_\mu \right\} d\psi \right]$$
(A2.6)

simplifying

$$\int \left((\kappa_r \theta_r - \sigma_r^2) \frac{2\kappa_r e^{-\kappa_r \psi}}{\sigma_r^2 (c_3 - e^{-\kappa_r \psi})} + (\kappa_\mu m(t) - \sigma_\mu^2) \frac{2\kappa_\mu e^{-\kappa_\mu \psi}}{\sigma_\mu^2 (d_3 - e^{-\kappa_\mu \psi})} - \kappa_r - \kappa_\mu \right) d\psi$$

$$= \frac{2(\kappa_r \theta_r - \sigma_r^2)}{\sigma_r^2} \ln|c_3 - e^{-\kappa_r \psi}| + \frac{2(\kappa_\mu m(t) - \sigma_\mu^2)}{\sigma_\mu^2} \ln|d_3 - e^{-\kappa_\mu \psi}| - \{\kappa_r + \kappa_\mu\}\psi \qquad (A2.7)$$

The integrating factor can be represented as

$$R(\psi) = \left|c_3 - e^{-\kappa_r \psi}\right|^{\frac{2(\kappa_r \theta_r - \sigma_r^2)}{\sigma_r^2}} \times \left|d_3 - e^{-\kappa_\mu \psi}\right|^{\frac{2(\kappa_\mu m(t) - \sigma_\mu^2)}{\sigma_\mu^2}} \times \exp\left[-\left\{\kappa_r + \kappa_\mu\right\}\psi\right]$$
(A2.8)

Simplify some of the terms, we let

$$\Phi_r = \kappa_r \theta_r$$
$$\Phi_\mu = \kappa_\mu m(t)$$

and rewrite equation (A2.8) as

$$R(\psi) = \left|c_3 - e^{-\kappa_r \psi}\right|^{\frac{2(\Phi_r - \sigma_R^2)}{\sigma_r^2}} \times \left|d_3 - e^{-\kappa_\mu \psi}\right|^{\frac{2(\Phi_\mu - \sigma_\mu^2)}{\sigma_\mu^2}} \times \exp\left[-\left\{\kappa_r + \kappa_\mu\right\}\psi\right]$$
(A2.9)

We can solve equation (A2.5) by writing it as

$$\frac{d}{d\psi} \Big(R(\psi) \tilde{G} \Big) = R(\psi) \Big[f_1(\psi, s_\mu) + f_2(\psi, s_r) \Big]$$
(A2.10)

integrating gives

$$R(\psi)\tilde{G} = \int_0^{\psi} R(t) \Big[f_1(t, s_{\mu}) + f_2(t, s_r) \Big] dt + c_4$$

then

$$\tilde{G} = \frac{1}{R(\psi)} \left\{ \int_0^{\psi} R(t) \Big[f_1(t, s_{\mu}) + f_2(t, s_r) \Big] dt + c_4 \right\}$$

This can be represented as

$$\tilde{G} = |c_3 - e^{-\kappa_r \psi}| \frac{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}}{\sigma_r^2} \times |d_3 - e^{-\kappa_\mu \psi}| \frac{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}}{\sigma_\mu^2} \times \exp\left[\{\kappa_r + \kappa_\mu\}\psi\right] \times \left\{\int_0^{\psi} \left[f_1(t, s_\mu) + f_2(t, s_r)\right] \times |c_3 - e^{-\kappa_r t}| \frac{\frac{2(\Phi_r - \sigma_r^2)}{\sigma_r^2}}{\sigma_r^2} \times |d_3 - e^{-\kappa_\mu t}| \frac{\frac{2(\Phi_\mu - \sigma_\mu^2)}{\sigma_\mu^2}}{\sigma_\mu^2} \times \exp\left[-\{\kappa_r + \kappa_\mu\}t\right] dt + c_4\right\}$$
(A2.11)

Here c_4 is a constant and can be determined when $\psi = 0$, equation (A2.11) can then be expressed as

$$\tilde{G}(0, s_r, s_\mu) = \left| c_3 - 1 \right|^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \times \left| d_3 - 1 \right|^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \times c_4 \tag{A2.12}$$

Rearranging in terms of c_4 and writing c_4 as a function of c_3 and d_3 we have

$$A(c_3, d_3) = \left| c_3 - 1 \right|^{\frac{2(\Phi_r - \sigma_r^2)}{\sigma_r^2}} \times \left| d_3 - 1 \right|^{\frac{2(\Phi_\mu - \sigma_\mu^2)}{\sigma_\mu^2}} \times \tilde{G}\left(0, \frac{-2\kappa_r}{\sigma_r^2(c_3 - 1)}, \frac{-2\kappa_\mu}{\sigma_\mu^2(d_3 - 1)}\right)$$
(A2.13)

In terms of the constant from equation (A2.11) we can write

$$\begin{aligned} |c_{3} - e^{-\kappa_{r}\psi}| \frac{2(\sigma_{r}^{2} - \Phi_{r})}{\sigma_{r}^{2}} \times |d_{3} - e^{-\kappa_{\mu}\psi}| \frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}} \times \exp\left[\{\kappa_{r} + \kappa_{\mu}\}\psi\right] \times A(c_{3}, d_{3}) \\ = \left|\frac{c_{3} - e^{-\kappa_{r}\psi}}{c_{3} - 1}\right|^{\frac{2(\sigma_{r}^{2} - \Phi_{r})}{\sigma_{r}^{2}}} \times \left|\frac{d_{3} - e^{-\kappa_{\mu}\psi}}{d_{3} - 1}\right|^{\frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r} + \kappa_{\mu}\}\psi\right] \times \tilde{G}\left(0, \frac{-2\kappa_{r}}{\sigma_{r}^{2}(c_{3} - 1)}, \frac{-2\kappa_{\mu}}{\sigma_{\mu}^{2}(d_{3} - 1)}\right)\right) \\ = \left|\frac{2\kappa_{r}e^{-\kappa_{r}\psi}}{\sigma_{r}^{2}s_{r}(1 - e^{-\kappa_{r}\psi}) + 2\kappa_{r}e^{-\kappa_{r}\psi}}\right|^{\frac{2(\sigma_{r}^{2} - \Phi_{r})}{\sigma_{r}^{2}}} \times \left|\frac{2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}{\sigma_{\mu}^{2}s_{\mu}(1 - e^{-\kappa_{\mu}\psi}) + 2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right|^{\frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r} + \kappa_{\mu}\}\psi\right] \\ \times \tilde{G}\left(0, \frac{-2\kappa_{r}}{\sigma_{r}^{2}(c_{3} - 1)}, \frac{-2\kappa_{\mu}}{\sigma_{\mu}^{2}(d_{3} - 1)}\right)\right) \tag{A2.14}$$

Equation (A2.11) can be simplified to

$$\begin{split} \tilde{G} &= \left| \frac{2\kappa_r e^{-\kappa_r \psi}}{\sigma_r^2 s(1 - e^{-\kappa_r \psi}) + 2\kappa_r e^{-\kappa_r \psi}} \right|^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \times \left| \frac{2\kappa_\mu e^{-\kappa_\mu \psi}}{\sigma_\mu^2 s_\mu (1 - e^{-\kappa_\mu \psi}) + 2\kappa_\mu e^{-\kappa_\mu \psi}} \right|^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \times \exp\left[\{\kappa_r + \kappa_\mu\}\psi \right] \\ &\times \tilde{G}\left(0, \frac{-2\kappa_r}{\sigma_r^2 (c_3 - 1)}, \frac{-2\kappa_\mu}{\sigma_\mu^2 (d_3 - 1)} \right) + \left\{ \int_0^{\psi} \left[f_1(t, s_\mu) + f_2(t, s_r) \right] \times \left| \frac{c_3 - e^{-\kappa_r \psi}}{c_3 - e^{-\kappa_r t}} \right|^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \\ &\times \left| \frac{d_3 - e^{-\kappa_\mu \psi}}{d_3 - e^{-\kappa_\mu t}} \right|^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \times \exp\left[\{\kappa_r + \kappa_\mu\}(\psi - t) \right] dt \right\} \end{split}$$
(A2.15)

This reduces to

$$\begin{split} \tilde{G} &= \left(\frac{2\kappa_r e^{-\kappa_r \psi}}{\sigma_r^2 s_r (1 - e^{-\kappa_r \psi}) + 2\kappa_r e^{-\kappa_r \psi}}\right)^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \times \left(\frac{2\kappa_\mu e^{-\kappa_\mu \psi}}{\sigma_\mu^2 s_\mu (1 - e^{-\kappa_\mu \psi}) + 2\kappa_\mu e^{-\kappa_\mu \psi}}\right)^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \\ &\times \exp\left[\{\kappa_r + \kappa_\mu\}\psi\right] \times \tilde{G}\left(0, \frac{2s\kappa_r}{\sigma_r^2 s_r (1 - e^{-\kappa_r \psi}) + 2\kappa_r e^{-\kappa_r \psi}}, \frac{2s_\mu \kappa_\mu}{\sigma_\mu^2 s_\mu (1 - e^{-\kappa_\mu \psi}) + 2\kappa_\mu e^{-\kappa_\mu \psi}}\right)^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \\ &+ \left\{\int_0^\psi \left[f_1(t, s_\mu) + f_2(t, s_r)\right] \times \left(\frac{2\kappa_r e^{-\kappa_r \psi}}{\sigma_r^2 s(e^{-\kappa_r t} - e^{-\kappa_r \psi}) + 2\kappa_r e^{-\kappa_r \psi}}\right)^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \\ &\times \left(\frac{2\kappa_\mu e^{-\kappa_\mu \psi}}{\sigma_\mu^2 s_\mu (e^{-\kappa_\mu t} - e^{-\kappa_\mu \psi}) + 2\kappa_\mu e^{-\kappa_\mu \psi}}\right)^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \times \exp\left[\{\kappa_r + \kappa_\mu\}(\psi - t)\right]dt\right\} \quad (A2.16) \end{split}$$

We can solve $f_1(t, s_\mu)$ as $s_r \to \infty$ equation (A2.16) simplifies to

$$\begin{split} 0 =& \left(\frac{2\kappa_{r}e^{-\kappa_{r}\psi}}{\sigma_{r}^{2}(1-e^{-\kappa_{r}\psi})}\right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}{\sigma_{\mu}^{2}s_{\mu}(1-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}\psi\right] \\ & \times \tilde{G}\left(0,\frac{2\kappa_{r}}{\sigma_{r}^{2}(1-e^{-\kappa_{r}\psi})},\frac{2s_{\mu}\kappa_{\mu}}{\sigma_{\mu}^{2}s_{\mu}(1-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right) \\ & + \left\{\int_{0}^{\psi}\left[f_{1}(t,s_{\mu})\right] \times \left(\frac{2\kappa_{r}e^{-\kappa_{r}\psi}}{\sigma_{r}^{2}(e^{-\kappa_{r}t}-e^{-\kappa_{r}\psi})}\right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}{\sigma_{\mu}^{2}s_{\mu}(e^{-\kappa_{\mu}t}-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \\ & \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}(\psi-t)\right]dt\right\} \end{split}$$

This can be rearranged as

$$\begin{split} &-\tilde{G}\left(0,\frac{2\kappa_{r}}{\sigma_{r}^{2}(1-e^{-\kappa_{r}\psi})},\frac{2s_{\mu}\kappa_{\mu}}{\sigma_{\mu}^{2}s_{\mu}(1-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right) \\ &=\left\{\int_{0}^{\psi}f_{1}(t,s_{\mu})\times\left(\frac{1-e^{-\kappa_{r}\psi}}{e^{-\kappa_{r}t}-e^{-\kappa_{r}\psi}}\right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}}\times\left(\frac{\sigma_{\mu}^{2}s_{\mu}(1-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}{\sigma_{\mu}^{2}s_{\mu}(e^{-\kappa_{\mu}t}-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \\ &\times\exp\left[-\left\{\kappa_{r}+\kappa_{\mu}\right\}t\right]dt\right\} \end{split}$$

 let

$$\zeta_r^{-1} = 1 - e^{-\kappa_r t}, \qquad z_r^{-1} = 1 - e^{-\kappa_r \psi} \tag{A2.17}$$

$$\zeta_{\mu}^{-1} = 1 - e^{-\kappa_{\mu}t}, \qquad z_{\mu}^{-1} = 1 - e^{-\kappa_{\mu}\psi}$$
(A2.18)

define

$$g(\zeta_{r}) = f_{1}(t, s_{\mu}) \times \frac{\frac{\zeta_{r}^{2(\sigma_{r}^{2} - \Phi_{r})}}{\sigma_{r}^{2}}}{\zeta_{r}(\zeta_{r})} \times \left(\frac{\zeta_{\mu}[\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)]}{\sigma_{\mu}^{2}s_{\mu}(\zeta_{\mu} - z_{1}) + 2\kappa_{\mu}\zeta_{\mu}(z_{\mu} - 1)}\right)^{\frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}}} \exp\left[-\{\kappa_{r} + \kappa_{\mu}\}t\right]$$
(A2.19)

substituting with a change of integral variable and limits

$$\int_{z_r}^{\infty} g(\zeta_r)(\zeta_r - z_r)^{-\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} d\zeta_r = -\kappa_r \tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu \alpha_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\right)$$

By the definition of Laplace transform

$$\tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\right) = \int_0^\infty \int_0^\infty \exp\left\{-\frac{2\kappa_r z_r}{\sigma_r^2}r - \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\mu\right\} G(0, r, \mu) dr d\mu$$
(A2.20)

We introduce the Gamma function such that

$$\begin{split} \tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu\kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu(z_\mu - 1)}\right) \\ &= \frac{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \times \int_0^\infty \int_0^\infty \exp\left\{-\frac{2\kappa_r z_r}{\sigma_r^2} r - \frac{2s_\mu\kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu(z_\mu - 1)}\mu\right\} \times G(0, r, \mu) dr d\mu \end{split}$$
(A2.21)

expanding the Gamma function

$$\tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\right) = \frac{1}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \times \int_0^\infty \int_0^\infty \int_0^\infty e^{-a} a^{\frac{2\Phi_r}{\sigma_r^2} - 2} \times \exp\left\{-\frac{2\kappa_r z_r}{\sigma_r^2} r - \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\mu\right\} \times G(0, r, \mu) dadr d\mu$$
(A2.22)

Make the substitution $a = \left(\frac{2\kappa_r r}{\sigma_r^2}\right)y$ we obtain

$$\begin{split} \tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\right) &= \frac{1}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \times \int_0^\infty \int_0^\infty \int_0^\infty \exp\left\{-\frac{2\kappa_r r}{\sigma_r^2} y\right\} \\ \left(\frac{2\kappa_r r}{\sigma_r^2} y\right)^{\frac{2\Phi_r}{\sigma_r^2} - 2} \times \left(\frac{2\kappa_r r}{\sigma_r^2}\right) \times \exp\left\{-\frac{2\kappa_r z_r}{\sigma_r^2} r - \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\mu\right\} \times G(0, r, \mu) dy dr d\mu \end{split}$$

$$(A2.23)$$

this can be rearranged to

$$\begin{split} \tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu \alpha_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\right) &= \frac{1}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \times \int_0^\infty \int_0^\infty G(0, r, \mu) \left(\frac{2\kappa_r r}{\sigma_r^2}\right)^{\frac{2\Phi_r}{\sigma_r^2} - 1} \\ &\times \exp\left\{-\frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\mu\right\} \left[\int_0^\infty \exp\left\{-\left(\frac{2\kappa_r r}{\sigma_r^2}\right)(y + z_r)\right\} y^{\left(\frac{2\Phi_r}{\sigma_r^2} - 2\right)} dy\right] dr d\mu \\ &\qquad (A2.24) \end{split}$$

let $\rho = y + z_r$ and substituting into (A2.24)

$$\begin{split} \tilde{G}\left(0, \frac{2\kappa_r z_r}{\sigma_r^2}, \frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\right) &= \frac{1}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \times \int_0^\infty \int_0^\infty G(0, r, \mu) \left(\frac{2\kappa_r r}{\sigma_r^2}\right)^{\frac{2\Phi_r}{\sigma_r^2} - 1} \\ &\times \exp\left\{-\frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\mu\right\} \times \left[\int_{z_r}^\infty \exp\left\{-\left(\frac{2\kappa_r r}{\sigma^2}\right)\varrho\right\} (\varrho - z_r)^{-\frac{2(\sigma_r^2 - \Phi_r)}{\sigma^2}} d\varrho\right] dr d\mu \\ &\qquad (A2.25) \end{split}$$

rearranging gives

$$\tilde{G}\left(0,\frac{2\kappa_{r}z_{r}}{\sigma_{r}^{2}},\frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)}\right) = \int_{z_{r}}^{\infty}\left(\varrho-z_{r}\right)^{-\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left[\int_{0}^{\infty}\int_{0}^{\infty}\frac{G(0,r,\mu)}{\Gamma(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1)} \times \left(\frac{2\kappa_{r}r}{\sigma_{r}^{2}}\right)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1} \exp\left\{-\frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)}\mu\right\} \exp\left\{-\left(\frac{2\kappa_{r}\varrho}{\sigma_{r}^{2}}\right)r\right\}drd\mu\right]d\varrho \quad (A2.26)$$

Thus we have shown that

$$g(\zeta_r) = -\kappa_r \int_0^\infty \int_0^\infty \frac{G(0,r,\mu)}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \left(\frac{2\kappa_r r}{\sigma_r^2}\right)^{\frac{2\Phi_r}{\sigma_r^2} - 1} \times \exp\left\{-\frac{2s_\mu \kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu (z_\mu - 1)}\mu\right\} \exp\left\{-\left(\frac{2\kappa_r \varrho}{\sigma_r^2}\right)r\right\} dr d\mu$$
(A2.27)

the initial condition from (3.4)

$$G(0, r, \mu; r_0, \mu_0) = \delta(r - r_0)\delta(\mu - \mu_0),$$

substituting into (A2.27)

$$g(\zeta_r) = -\kappa_r \int_0^\infty \int_0^\infty \frac{\delta(r-r_0)\delta(\mu-\mu_0)}{\Gamma(\frac{2\Phi_r}{\sigma_r^2}-1)} \left(\frac{2\kappa_r r}{\sigma_r^2}\right)^{\frac{2\Phi_r}{\sigma_r^2}-1} \times \exp\left\{-\frac{2s_\mu\kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\alpha_\mu(z_\mu-1)}\mu\right\} \exp\left\{-\left(\frac{2\kappa_r \varrho}{\sigma_r^2}\right)r\right\} dr d\mu$$
(A2.28)

using the properties of the Dirac Delta function

$$g(\zeta_r) = \frac{-\kappa_r}{\Gamma(\frac{2\Phi_r}{\sigma_r^2} - 1)} \left(\frac{2\kappa_r r_0}{\sigma_r^2}\right)^{\frac{2\Phi_r}{\sigma_r^2} - 1} \times \exp\left\{-\frac{2s_\mu\kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu(z_\mu - 1)}\mu\right\} \exp\left\{-\left(\frac{2\kappa_r \zeta_r}{\sigma_r^2}\right)r_0\right\}$$
(A2.29)

We can then substitute can into (A2.19) to get an explicit form of $f_1(t,s_\mu)$

$$\begin{split} f_{1}(t,s_{\mu}) = & \frac{-\kappa_{r}}{\Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)} \left(\frac{2\kappa_{r}r_{0}}{\sigma_{r}^{2}}\right)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1} \times \exp\left\{-\frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)}\mu_{0}\right\} \times \exp\left\{-\left(\frac{2\kappa_{r}\zeta_{r}}{\sigma_{r}^{2}}\right)r_{0}\right\} \\ \times & \frac{\zeta_{r}(\zeta_{r}-1)}{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\zeta_{r}}} \times \left(\frac{\zeta_{\mu}[\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)]}{\sigma_{\mu}^{2}s_{\mu}(\zeta_{\mu}-z_{\mu})+2\kappa_{\mu}\zeta_{\mu}(z_{\mu}-1)}\right)^{\frac{2(\Phi_{\mu}-\sigma_{\mu}^{2})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}t\right] \end{split}$$

$$(A2.30)$$

By a similar operation we can show

$$f_{2}(t,s_{r}) = \frac{-\kappa_{\mu}}{\Gamma(\frac{2\phi_{\mu}}{\sigma_{\mu}^{2}}-1)} \left(\frac{2\kappa_{\mu}\mu_{0}}{\sigma_{\mu}^{2}}\right)^{\frac{2\phi_{\mu}}{\sigma_{\mu}^{2}}-1} \times \exp\left\{-\frac{2s_{r}\kappa_{r}z_{r}}{\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)}r_{0}\right\} \times \exp\left\{-\left(\frac{2\kappa_{\mu}\zeta_{\mu}}{\sigma_{\mu}^{2}}\right)\mu_{0}\right\}$$
$$\times \frac{\zeta_{\mu}(\zeta_{\mu}-1)}{\frac{2(\sigma_{\mu}^{2}-\phi_{\mu})}{\zeta_{\mu}}} \times \left(\frac{\zeta_{r}[\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)]}{\sigma_{r}^{2}s_{r}(\zeta_{r}-z_{r})+2\kappa_{r}\zeta_{r}(z_{r}-1)}\right)^{\frac{2(\phi_{r}-\sigma_{r}^{2})}{\sigma_{r}^{2}}} \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}t\right]$$
(A2.31)

The first component on the RHS of equantion (A2.16) can be written as

$$J_{1} = \left(\frac{2\kappa_{r}e^{-\kappa_{r}\psi}}{\sigma_{r}^{2}s_{r}(1-e^{-\kappa_{r}\psi})+2\kappa_{r}e^{-\kappa_{r}\psi}}\right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}{\sigma_{\mu}^{2}s_{\mu}(1-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}\psi\right] \times \tilde{G}\left(0, \frac{2s_{r}\kappa_{r}}{\sigma_{r}^{2}s_{r}(1-e^{-\kappa_{r}\psi})+2\kappa_{r}e^{-\kappa_{r}\psi}}, \frac{2s_{\mu}\kappa_{\mu}}{\sigma_{\mu}^{2}s_{\mu}(1-e^{-\kappa_{\mu}\psi})+2\kappa_{\mu}e^{-\kappa_{\mu}\psi}}\right)^{(A2.32)}$$

Making use of equation (A2.17), we can write

$$J_{1} = \left(\frac{2\kappa_{r}(z_{r}-1)}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r}-1)}\right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\alpha_{\mu}(z_{\mu}-1)}{\sigma_{\mu}^{2}s + 2\kappa_{\mu}(z_{\mu}-1)}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \exp\left[\{\kappa_{r}+\kappa_{\mu}\}\psi\right] \times \tilde{G}\left(0, \frac{2s_{r}\kappa_{r}z_{r}}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r}-1)}, \frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu}-1)}\right)$$
(A2.33)

Substituting the initial conditions equation (3.10) yields

$$J_{1} = \left(\frac{2\kappa_{r}(z_{r}-1)}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r}-1)}\right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\kappa_{\mu}(z_{\mu}-1)}{\sigma_{\mu}^{2}s + 2\kappa_{\mu}(z_{\mu}-1)}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \exp\left[\{\kappa_{r}+\kappa_{\mu}\}\psi\right] \times \exp\left\{-\frac{2s_{r}\kappa_{r}z_{r}}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r}-1)}r_{0} - \frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu}-1)}\mu_{0}\right\}$$
(A2.34)

The second component on the RHS of equation (A2.16) can be written as

$$J_{2} = \int_{z_{r}}^{\infty} \left[f_{1}(t,s_{\mu}) \right] \times \left(\frac{2\kappa_{r}\zeta_{r}(z_{r}-1)}{\sigma_{r}^{2}s(\zeta_{r}-z_{r})+2\kappa_{r}\zeta_{r}(z_{r}-1)} \right)^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\kappa_{\mu}\zeta_{\mu}(z_{\mu}-1)}{\sigma_{\mu}^{2}s_{\mu}(\zeta_{\mu}-z_{\mu})+2\kappa_{\mu}\zeta_{\mu}(z_{\mu}-1)} \right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}(\psi-t) \right] \frac{d\zeta_{r}}{\kappa_{r}\zeta_{r}(\zeta_{r}-1)}$$
(A2.35)

Substituting equation (A2.30) yields

$$J_{2} = \frac{-1}{\Gamma(\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1)} \int_{z_{r}}^{\infty} \left(\frac{2\kappa_{r}r_{0}}{\sigma_{r}^{2}}\right)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1} \times \exp\left\{-\frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)}\mu_{0}\right\} \times \exp\left\{-\left(\frac{2\kappa_{r}\zeta_{r}}{\sigma_{r}^{2}}\right)r_{0}\right\}$$
$$\times \left(\frac{2\kappa_{\mu}(z_{\mu} - 1)}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)}\right)^{\frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \left(\frac{2\kappa_{r}(z_{r} - 1)}{\sigma_{r}^{2}s(\zeta_{r} - z_{r}) + 2\kappa_{r}\zeta_{r}(z_{r} - 1)}\right)^{\frac{2(\sigma_{r}^{2} - \Phi_{r})}{\sigma_{r}^{2}}} \times \exp\left[\{\kappa_{r} + \kappa_{\mu}\}\psi\right]d\zeta_{r}$$
(A2.36)

Rearranging gives

$$J_{2} = \frac{-\left[2\kappa_{r}(z_{r}-1)\right]^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}}}{\Gamma(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1)} \left(\frac{2\kappa_{r}r_{0}}{\sigma_{r}^{2}}\right)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1} \times \exp\left\{-\frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)}\mu_{0}\right\} \times \left(\frac{2\kappa_{\mu}(z_{\mu}-1)}{\sigma_{\mu}^{2}s_{1}+2\kappa_{\mu}(z_{\mu}-1)}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\{\kappa_{r}+\kappa_{\mu}\}\psi\right]G_{1}(r_{0})$$
(A2.37)

where

$$G_1(r) = \int_{z_r}^{\infty} \exp\left\{-\left(\frac{2\kappa_r r}{\sigma_r^2}\right)\zeta_r\right\} \left[\sigma_r^2 s_r(\zeta_r - z_r) + 2\kappa_r \zeta_r(z_r - 1)\right]^{\frac{2(\Phi_r - \sigma_r^2)}{\sigma_r^2}} d\zeta_r \qquad (A2.38)$$

Let $y = \sigma_r^2 s(\zeta_r - z_r) + 2\kappa_r \zeta_r(z_r - 1)$, so that

$$d\zeta_r = \frac{dy}{\sigma_r^2 s_r + 2\kappa_r (z_r - 1)}$$
(A2.39)

substituting into equation (A2.38) we obtain

$$G_{1}(r) = \frac{1}{\sigma_{r}^{2}s_{r} + 2\kappa_{r}(z_{r}-1)} \exp\left\{-\left(\frac{2\kappa_{r}r}{\sigma_{r}^{2}}\right) \left(\frac{\sigma_{r}^{2}s_{r}z_{r}}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r}-1)}\right)\right\}$$
$$\int_{2\kappa_{r}z_{r}(z_{r}-1)}^{\infty} \exp\left\{-\left(\frac{2\kappa_{r}r}{\sigma_{r}^{2}}\right) \left(\frac{y}{\sigma_{r}^{2}s_{r} + 2\kappa_{r}(z_{r}-1)}\right)\right\} y^{\frac{2(\Phi_{r}-\sigma_{r}^{2})}{\sigma_{r}^{2}}} dy$$
(A2.40)

Now let

$$\xi = \left(\frac{2\kappa_r r}{\sigma_r^2}\right) \left(\frac{y}{\sigma_r^2 s + 2\kappa_r(z_r - 1)}\right) = \left(\frac{2\kappa_r r y}{\sigma_r^2 \left[\sigma_r^2 s + 2\kappa_r(z_r - 1)\right]}\right)$$
(A2.41)

This implies

$$dy = \frac{\sigma_r^2 [\sigma_r^2 s_r + 2\kappa_r (z_r - 1)]}{2\kappa_r r} d\xi$$
(A2.42)

substituting into equation (A2.40) yields

$$G_{1}(r) = \frac{\sigma_{r}^{2}}{2\kappa_{r}r} \exp\left\{-\left(\frac{2\kappa_{r}r}{\sigma_{r}^{2}}\right) \left(\frac{\sigma_{r}^{2}s_{r}z_{r}}{\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)}\right)\right\} \\ \left[\frac{\sigma_{r}^{2}[\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)]}{2\kappa_{r}r}\right]^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-2\right)} \int_{\frac{4\kappa_{r}^{2}rz_{r}(z_{r}-1)}{\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)}}^{\infty} \exp\left\{-\xi\right\} \xi^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)-1} d\xi \quad (A2.43)$$

rearranging and inserting the Gamma function

$$G_{1}(r) = \left(\sigma_{r}^{2}s + 2\kappa_{r}(z_{r}-1)\right)^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-2\right)} \left(\frac{\sigma_{r}^{2}}{2\kappa_{r}r}\right)^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)} \exp\left\{-\left(\frac{2\kappa_{r}r}{\sigma_{r}^{2}}\right)\left(\frac{\sigma_{r}^{2}s_{r}z_{r}}{\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)}\right)\right\}$$

$$\left[\Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)-\int_{0}^{\frac{4\kappa_{r}^{2}rz_{r}(z_{r}-1)}{\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)}}\exp\left\{-\xi\right\}\xi^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)-1}d\xi\right]$$
(A2.44)

substituting back into equation (A2.37) gives

$$J_{2} = \frac{-\left[2\kappa_{r}(z_{r}-1)\right]^{\frac{2(\sigma_{r}^{2}-\Phi_{r})}{\sigma_{r}^{2}}}}{\Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)} \left(\frac{2\kappa_{r}r_{0}}{\sigma_{r}^{2}}\right)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1} \times \exp\left\{-\frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)}\mu_{0}\right\}$$

$$\times \left(\frac{2\kappa_{\mu}(z_{\mu}-1)}{\sigma_{\mu}^{2}s_{\mu}+2\kappa_{\mu}(z_{\mu}-1)}\right)^{\frac{2(\sigma_{\mu}^{2}-\Phi_{\mu})}{\sigma_{\mu}^{2}}} \times \exp\left[\left\{\kappa_{r}+\kappa_{\mu}\right\}\psi\right] \left(\sigma_{r}^{2}s+2\kappa_{r}(z_{r}-1)\right)^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-2\right)} \left(\frac{\sigma_{r}^{2}}{2\kappa_{r}r_{0}}\right)^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)}$$

$$\exp\left\{-\left(\frac{2\kappa_{r}r_{0}}{\sigma_{r}^{2}}\right) \left(\frac{\sigma_{r}^{2}sz}{\sigma_{r}^{2}s+2\kappa_{r}(z_{r}-1)}\right)\right\} \left[\Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right) - \int_{0}^{\frac{4\kappa_{r}^{2}rz_{r}(z_{r}-1)}{\sigma_{r}^{2}s_{r}+2\kappa_{r}(z_{r}-1)}} \exp\left\{-\xi\right\}\xi^{\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right)-1}d\xi\right]$$

$$(A2.45)$$

this reduces to

$$J_{2} = \frac{-1}{\Gamma(\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1)} \times \exp\left\{-\frac{2\kappa_{r}s_{r}z_{r}}{\sigma_{r}^{2}s_{r} + 2\kappa_{r}(z_{r} - 1)}r_{0} - \frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)}\mu_{0}\right\}$$
$$\times \left(\frac{2\kappa_{r}(z_{r} - 1)}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r} - 1)}\right)^{\frac{2(\sigma_{r}^{2} - \Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\kappa_{\mu}(z_{\mu} - 1)}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)}\right)^{\frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}}}$$
$$\times \exp\left[\left\{\kappa_{r} + \kappa_{\mu}\right\}\psi\right] \times \Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1\right)\left[1 - \Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1; \frac{4\kappa_{r}^{2}z_{r}(z_{r} - 1)}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r} - 1)}r_{0}\right)\right]$$
(A2.46)

Similarly the third component on the RHS of equation (A2.16) can be written as

$$J_{3} = \frac{-1}{\Gamma(\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 1)} \times \exp\left\{-\frac{2\kappa_{r}s_{r}z_{r}}{\sigma_{r}^{2}s_{r} + 2\kappa_{r}(z_{r} - 1)}r_{0} - \frac{2s_{\mu}\kappa_{\mu}z_{\mu}}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)}\mu_{0}\right\}$$
$$\times \left(\frac{2\kappa_{r}(z_{r} - 1)}{\sigma_{r}^{2}s + 2\kappa_{r}(z_{r} - 1)}\right)^{\frac{2(\sigma_{r}^{2} - \Phi_{r})}{\sigma_{r}^{2}}} \times \left(\frac{2\alpha_{\mu}(z_{\mu} - 1)}{\sigma_{\mu}^{2}s_{\mu} + 2\kappa_{\mu}(z_{\mu} - 1)}\right)^{\frac{2(\sigma_{\mu}^{2} - \Phi_{\mu})}{\sigma_{\mu}^{2}}}$$
$$\times \exp\left[\{\kappa_{r} + \kappa_{\mu}\}\psi\right] \times \Gamma\left(\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 1\right)\left[1 - \Gamma\left(\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 1; \frac{4\alpha_{\mu}^{2}z_{\mu}(z_{\mu} - 1)}{\sigma_{\mu}^{2}s + 2\kappa_{\mu}(z_{\mu} - 1)}\mu_{0}\right)\right]$$
(A2.47)

By combining J1, J2 and J3 equation (A2.16) becomes

$$\begin{split} \tilde{G} &= \left(\frac{2\kappa_r(z_r - 1)}{\sigma_r^2 s + 2\kappa_r(z_r - 1)}\right)^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \times \left(\frac{2\kappa_\mu(z_\mu - 1)}{\sigma_\mu^2 s_\mu + 2\kappa_\mu(z_\mu - 1)}\right)^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \\ &= \exp\left[\{\kappa_r + \kappa_\mu\}\psi\right] \times \exp\left\{-\frac{2s_r\kappa_r z_r}{\sigma_r^2 s_r + 2\kappa_r(z_r - 1)}r_0 - \frac{2s_\mu\kappa_\mu z_\mu}{\sigma_\mu^2 s_\mu + 2\kappa_\mu(z_\mu - 1)}\mu_0\right\} \\ &\times \left[\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1; \frac{4\kappa_r^2 z_r(z_r - 1)}{\sigma_r^2 s_r + 2\kappa_r(z_r - 1)}r_0\right) + \Gamma\left(\frac{2\Phi_\mu}{\sigma_\mu^2} - 1; \frac{4\kappa_\mu^2 z_\mu(z_\mu - 1)}{\sigma_\mu^2 s_\mu + 2\kappa_\mu(z_\mu - 1)}\mu_0\right) - 1\right] \\ &\quad (A2.48) \end{split}$$

and finally

$$\begin{split} \tilde{G} &= \left(\frac{2\kappa_r}{\sigma_r^2 s(e^{\kappa_r\psi} - 1) + 2\kappa_r}\right)^{\frac{2(\sigma_r^2 - \Phi_r)}{\sigma_r^2}} \times \left(\frac{2\kappa_\mu}{\sigma_\mu^2 s_\mu(e^{\kappa_\mu\psi} - 1) + 2\kappa_\mu}\right)^{\frac{2(\sigma_\mu^2 - \Phi_\mu)}{\sigma_\mu^2}} \\ &= \exp\left[\{\kappa_r + \kappa_\mu\}\psi\right] \times \exp\left\{-\frac{2s_r\kappa_r e^{\kappa_r\psi}}{\sigma_r^2 s(e^{\kappa_r\psi} - 1) + 2\kappa_r}r_0 - \frac{2s_\mu\kappa_\mu e^{\kappa_\mu\psi}}{\sigma_\mu^2 s_\mu(e^{\kappa_\mu\psi} - 1) + 2\kappa_\mu}\mu_0\right\} \\ &\times \left[\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1; \frac{4\kappa_r^2 e^{\kappa_r\psi}}{(e^{\kappa_r\psi} - 1)(\sigma_r^2 s(e^{\kappa_r\psi} - 1) + 2\kappa_r)}r_0\right) \\ &+ \Gamma\left(\frac{2\Phi_\mu}{\sigma_\mu^2} - 1; \frac{4\kappa_\mu^2 e^{\kappa_\mu\psi}}{(e^{\kappa_\mu\psi} - 1)(\sigma_\mu^2 s_\mu(e^{\kappa_\mu\psi} - 1) + 2\kappa_\mu)}\mu_0\right) - 1\right] \end{split}$$
(A2.49)

Appendix 3. Proof of Proposition 3.4

Start by making the following transforms

$$A_{r} = \frac{2\kappa_{r}r_{0}}{\sigma_{r}^{2}(1 - e^{-\kappa_{r}\psi})}$$

$$A_{\mu} = \frac{2\kappa_{\mu}\mu_{0}}{\sigma_{\mu}^{2}(1 - e^{-\kappa_{\mu}\psi})}$$

$$z_{r} = \frac{\sigma_{r}^{2}s(e^{\kappa_{r}\psi} - 1) + 2\kappa_{r}}{2\kappa_{r}}$$

$$z_{\mu} = \frac{\sigma_{\mu}^{2}s_{\mu}(e^{\kappa_{\mu}\psi} - 1) + 2\kappa_{\mu}}{2\kappa_{\mu}}$$
(A3.1)

 $\quad \text{and} \quad$

$$h = \exp\left[\{\kappa_r + \kappa_\mu\}\psi\right] \tag{A3.2}$$

We can rewrite equation (A2.49) as

$$\begin{split} \tilde{G} = & h \times z_r^{\frac{2\Phi_r}{\sigma_r^2} - 2} \times z_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^2} - 2} \exp\left\{-\frac{A_r}{z_r}(z_r - 1)\right\} \times \exp\left\{-\frac{A_{\mu}}{z_{\mu}}(z_{\mu} - 1)\right\} \\ & \times \left[\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1; \frac{A_r}{z_r}\right) + \Gamma\left(\frac{2\Phi_{\mu}}{\sigma_{\mu}^2} - 1; \frac{A_{\mu}}{z_{\mu}}\right) - 1\right] \\ = & h \times z_r^{\frac{2\Phi_r}{\sigma_r^2} - 2} \times z_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^2} - 2} \exp\left\{-\frac{A_r}{z_r}(z_r - 1)\right\} \times \exp\left\{-\frac{A_{\mu}}{z_{\mu}}(z_{\mu} - 1)\right\} \\ & \times \left[\frac{1}{\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1\right)} \int_0^{\frac{A_r}{z_r}} e^{-\beta_r} \beta_r^{\frac{2\Phi_r}{\sigma_r^2} - 2} d\beta_r + \frac{1}{\Gamma\left(\frac{2\Phi_{\mu}}{\sigma_{\mu}^2} - 1\right)} \int_0^{\frac{A_{\mu}}{z_{\mu}}} e^{-\beta_{\mu}} \beta_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^2} - 2} d\beta_{\mu} - 1\right] \end{split}$$

Separate \tilde{G} into 3 parts

$$\begin{split} \tilde{F}_{1} = h \times z_{r}^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 2} \times z_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 2} \exp\left\{-\frac{A_{r}}{z_{r}}(z_{r}-1)\right\} \times \exp\left\{-\frac{A_{\mu}}{z_{\mu}}(z_{\mu}-1)\right\} \times \frac{1}{\Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1\right)} \int_{0}^{\frac{A_{r}}{z_{r}}} e^{-\beta_{r}} \beta_{r}^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 2} d\beta_{r} \\ (A3.3) \\ \tilde{F}_{2} = h \times z_{r}^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 2} \times z_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 2} \exp\left\{-\frac{A_{r}}{z_{r}}(z_{r}-1)\right\} \times \exp\left\{-\frac{A_{\mu}}{z_{\mu}}(z_{\mu}-1)\right\} \times \frac{1}{\Gamma\left(\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 1\right)} \int_{0}^{\frac{A_{\mu}}{z_{\mu}}} e^{-\beta_{\mu}} \beta_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 2} d\beta_{\mu} \\ (A3.4) \\ \tilde{F}_{3} = -h \times z_{r}^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 2} \times z_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 2} \exp\left\{-\frac{A_{r}}{z_{r}}(z_{r}-1)\right\} \times \exp\left\{-\frac{A_{\mu}}{z_{\mu}}(z_{\mu}-1)\right\} \tag{A3.5}$$

Starting with \tilde{F}_1 , let $\xi = 1 - \frac{z_r}{A_r} \beta_r$ then

$$\tilde{F}_{1} = h \times e^{-(A_{r} + A_{\mu})} \times \frac{A_{r}^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1}}{\Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 1\right)} e^{\frac{A_{\mu}}{z_{\mu}}} z_{\mu}^{\frac{2\Phi_{\mu}}{\sigma_{\mu}^{2}} - 2} \times \int_{0}^{e^{\frac{A_{r}}{z_{r}}}} (1 - \xi)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}} - 2} z_{r}^{-1} e^{-\frac{A_{r}}{z_{r}}\xi} d\xi \qquad (A3.6)$$

from equation (A3.1) we can express the variable s_r and s_μ as

$$s_r = \frac{2\kappa_r(z_r - 1)}{\sigma_r^2(e^{\kappa_r\psi} - 1)}, \quad s_\mu = \frac{2\alpha_\mu(z_\mu - 1)}{\sigma_\mu^2(e^{\kappa_\mu\psi} - 1)}$$

representing the Laplace transform as

$$\mathcal{L}\{\hat{F}_1\} = \int_0^\infty \int_0^\infty \exp\left\{-\frac{2\kappa_r(z_r-1)}{\sigma_r^2(e^{\kappa_r\psi}-1)}r\right\} \times \exp\left\{-\frac{2\kappa_\mu(z_\mu-1)}{\sigma_\mu^2(e^{\kappa_\mu\psi}-1)}\mu\right\} \hat{F}_1 dr d\mu$$

substituting

$$y_r = \frac{2\kappa_r r}{\sigma_r^2 (e^{\kappa_r \psi} - 1)}, \quad y_\mu = \frac{2\kappa_\mu \mu}{\sigma_\mu^2 (e^{\kappa_\mu \psi} - 1)}$$
(A3.7)

then

$$\begin{aligned} \mathcal{L}\{\hat{F}_{1}\} = & \frac{\sigma_{r}^{2}(e^{\kappa_{r}\psi}-1)}{2\kappa_{r}} \frac{\sigma_{\mu}^{2}(e^{\kappa_{\mu}\psi}-1)}{2\kappa_{\mu}} \int_{0}^{\infty} \int_{0}^{\infty} e^{-y_{r}z_{r}} e^{-y_{\mu}z_{\mu}} \times e^{y_{r}} e^{y_{\mu}} \hat{F}_{1} dy_{r} \\ = & \frac{\sigma_{r}^{2}(e^{\kappa_{r}\psi}-1)}{2\kappa_{r}} \frac{\sigma_{\mu}^{2}(e^{\kappa_{\mu}\psi}-1)}{2\kappa_{\mu}} \mathcal{L}\{e^{y_{r}}e^{y_{\mu}}\hat{H}\} \end{aligned}$$

and in terms of the inverse Laplace transform

$$\mathcal{L}^{-1}\{\tilde{F}_1\} = \frac{2\kappa_r}{\sigma_r^2(e^{\kappa_r\psi} - 1)} \frac{2\kappa_\mu}{\sigma_\mu^2(e^{\kappa_\mu\psi} - 1)} e^{-y_r} e^{-y_\mu} \mathcal{L}^{-1}\{\tilde{F}_1\}$$

Applying the inverse transform to (A3.6) we obtain

$$\begin{split} \hat{F}_1 = & h \times e^{-(A_r + A_\mu)} \times \frac{A_r^{\frac{2\Phi_r}{\sigma_r^2} - 1}}{\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1\right)} \times \frac{2\kappa_r}{\sigma_r^2(e^{\kappa_r\psi} - 1)} \frac{2\kappa_\mu}{\sigma_\mu^2(e^{\kappa_\mu\psi} - 1)} e^{-y_r} e^{-y_\mu} \\ & \times \int_0^{e^{\frac{A_r}{z_r}}} (1 - \xi)^{\frac{2\Phi_r}{\sigma_r^2} - 2} \mathcal{L}^{-1} \left\{ e^{\frac{A_\mu}{z_\mu}} z_\mu^{\frac{2\Phi_\mu}{\sigma_\mu^2} - 2} \times e^{\frac{A_r}{z_r}\xi} z^{-1} \right\} d\xi \end{split}$$

The inverse Laplace can be solved as

$$\begin{split} \hat{F}_1 = & h \times e^{-(A_r + A_\mu)} \times \frac{A_r^{\frac{2\Phi_r}{\sigma_r^2} - 1}}{\Gamma\left(\frac{2\Phi_r}{\sigma_r^2} - 1\right)} \times \frac{2\kappa_r}{\sigma_r^2 (e^{\kappa_r \psi} - 1)} \frac{2\kappa_\mu}{\sigma_\mu^2 (e^{\kappa_\mu \psi} - 1)} e^{-y_r} e^{-y_\mu} \\ & \times \left(\frac{y_\mu}{A_\mu}\right)^{\frac{1}{2} - \frac{\Phi_\mu}{\sigma_\mu^2}} I_{1 - \frac{2\Phi_\mu}{\sigma_\mu^2}} (2\sqrt{A_\mu y_\mu}) \int_0^1 (1 - \xi)^{\frac{2\Phi_r}{\sigma_r^2} - 2} I_0 (2\sqrt{A_r y_r}\xi) d\xi \end{split}$$

We can simplify

$$\int_{0}^{1} (1-\xi)^{\frac{2\Phi_{r}}{\sigma_{r}^{2}}-2} I_{0}(2\sqrt{A_{r}y_{r}\xi})d\xi = \Gamma\left(\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1\right) (A_{r}y_{r})^{\frac{1}{2}-\frac{\Phi_{r}}{\sigma_{r}^{2}}} I_{\frac{2\Phi_{r}}{\sigma_{r}^{2}}-1}(2\sqrt{A_{r}y_{r}})$$

on substitution

$$\begin{split} \hat{F}_1 = & h \times e^{-(A_r + A_\mu)} \times \frac{2\kappa_r}{\sigma_r^2 (e^{\kappa_r \psi} - 1)} \frac{2\kappa_\mu}{\sigma_\mu^2 (e^{\kappa_\mu \psi} - 1)} e^{-y_r} e^{-y_\mu} \left(\frac{y_r}{A_r}\right)^{\frac{1}{2} - \frac{\Phi_r}{\sigma_r^2}} \times I_{1 - \frac{2\Phi_r}{\sigma_r^2}}(2\sqrt{A_r y_r}) \\ & \times \left(\frac{y_\mu}{A_\mu}\right)^{\frac{1}{2} - \frac{\Phi_\mu}{\sigma_\mu^2}} I_{1 - \frac{2\Phi_\mu}{\sigma_\mu^2}}(2\sqrt{A_\mu y_\mu}) \end{split}$$

similarly we can show $\hat{F}_2 = \hat{F}_1$ and $\hat{F}_3 = -\hat{F}_1$. This implies $\hat{G} = \hat{F}_1$. Substituting for (A3.1), (A3.2) and (A3.7) we obtain our result,

$$\hat{G} = \exp\left\{-\frac{2\kappa_r}{\sigma_r^2(e^{\kappa_r\psi}-1)}(r_0e^{\kappa_r\psi}+r) - \frac{2\kappa_\mu}{\sigma_\mu^2(e^{\kappa_\mu\psi}-1)}(\mu_0e^{\kappa_\mu\psi}+\mu)\right\}$$

$$\times \left(\frac{r_0e^{\kappa_r\psi}}{r}\right)^{\frac{\Phi_r}{\sigma_r^2}-\frac{1}{2}} \times \left(\frac{\mu_0e^{\kappa_\mu\psi}}{\mu}\right)^{\frac{\Phi_\mu}{\sigma_\mu^2}-\frac{1}{2}} I_{\frac{2\Phi_r}{\sigma_r^2}-1}\left(\frac{4\kappa_r}{\sigma_r^2(e^{\kappa_r\psi}-1)}(r\times r_0e^{\kappa_r\psi})^{\frac{1}{2}}\right)$$

$$\times \times I_{\frac{2\Phi_\mu}{\sigma_\mu^2}-1}\left(\frac{4\kappa_\mu}{\sigma_\mu^2(e^{\kappa_\mu\psi}-1)}(\mu\times\mu_0e^{\kappa_\mu\psi})^{\frac{1}{2}}\right) \times \frac{2\kappa_re^{\kappa_r\psi}}{\sigma_r^2(e^{\kappa_r\psi}-1)}\frac{2\kappa_\mu e^{\kappa_\mu\psi}}{\sigma_\mu^2(e^{\kappa_\mu\psi}-1)}$$
(A3.8)

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